

Simulation of 2D Maxwell-Bloch Equations

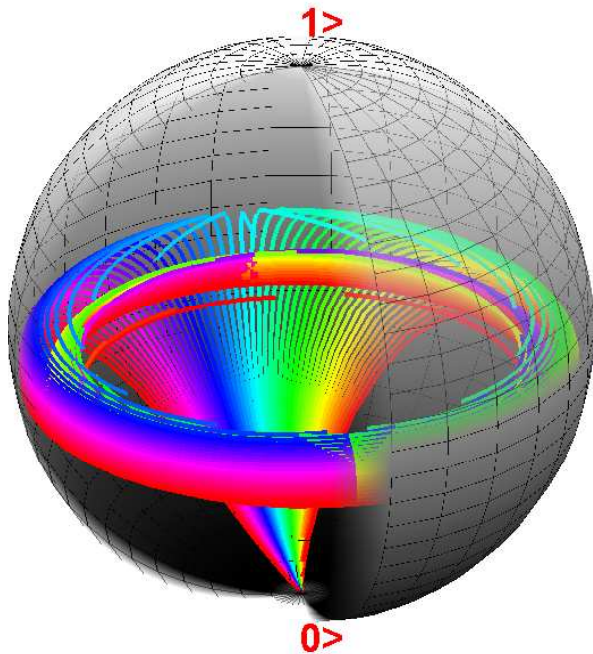
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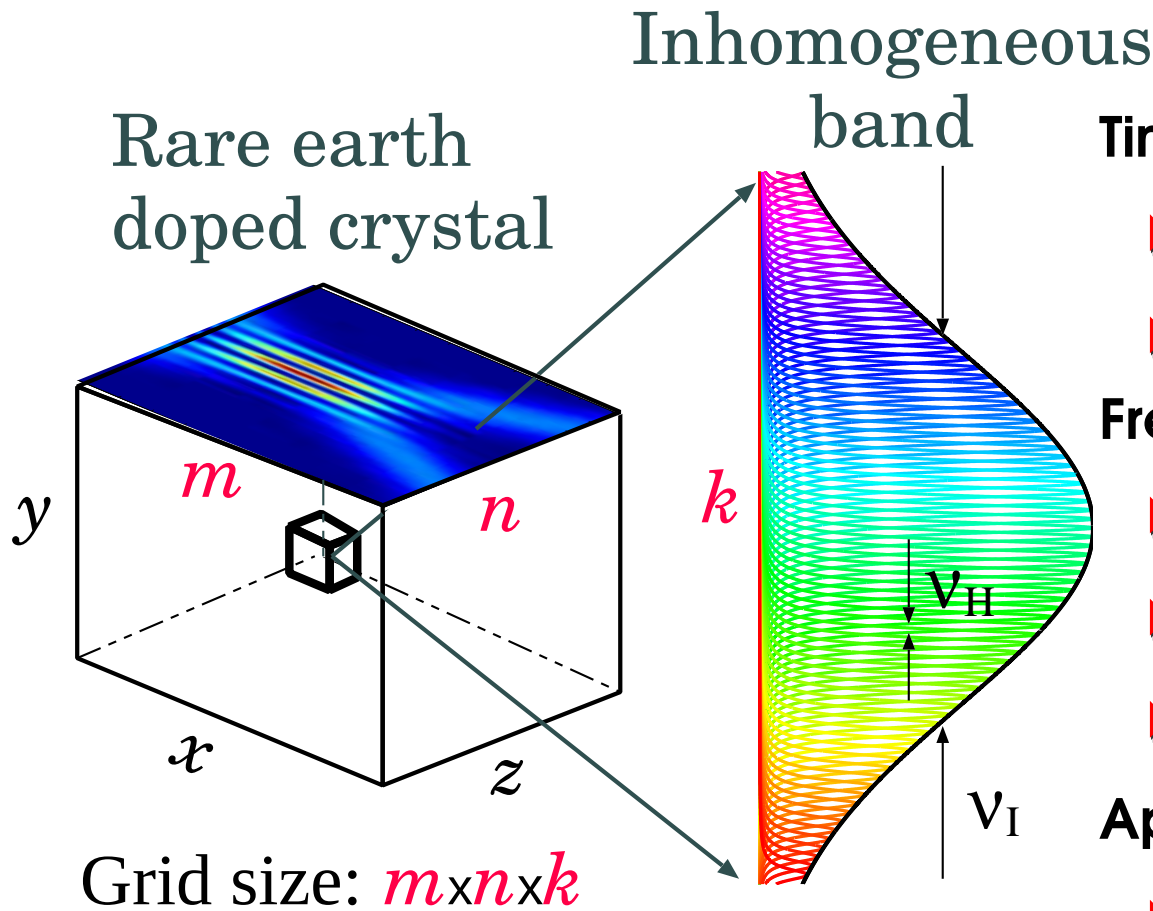
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-
- ▶ *Material property*
 - ▶ *Simulation schemes and results*
 - ▶ *Experimental results*

Inhomogeneously broadened material



Time domain parameters

- ▶ Life time $T_1 = 10\text{msec}$
- ▶ Decoherence time $T_2 = 5\text{ }\mu\text{sec}$

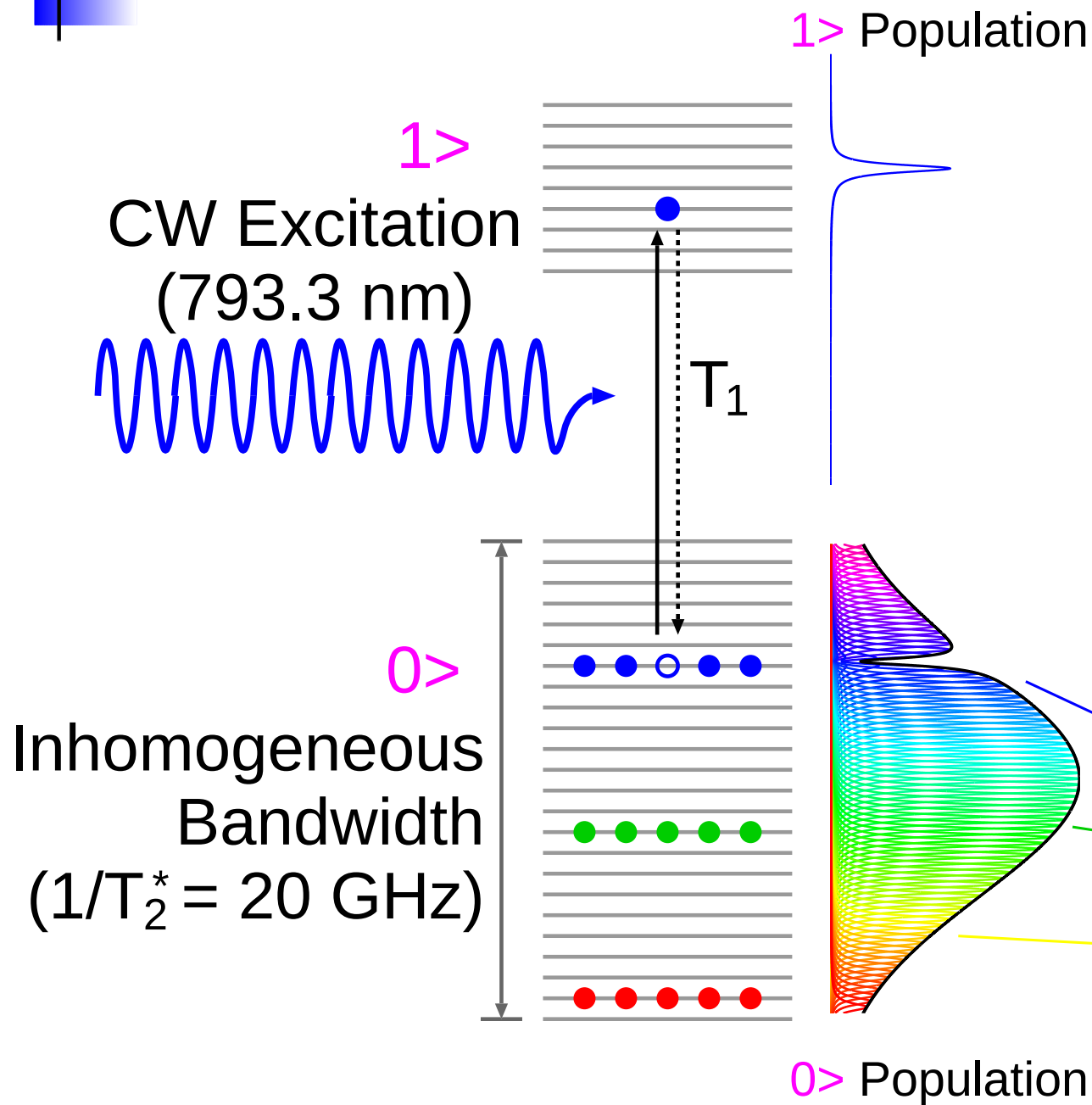
Frequency domain parameters

- ▶ Inhomogeneous BW $\nu_I = 20\text{GHz}$
- ▶ Homogeneous BW $\nu_H = 25\text{KHz}$
- ▶ Number of bins $N = \nu_I / \nu_H \approx 10^6$

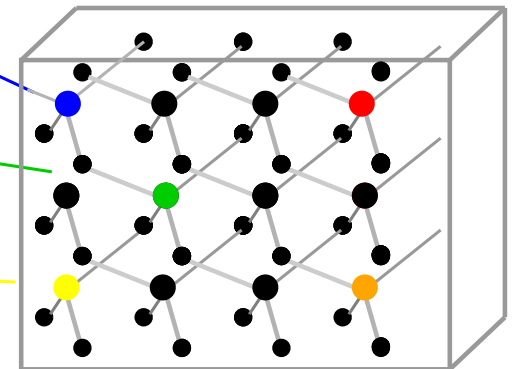
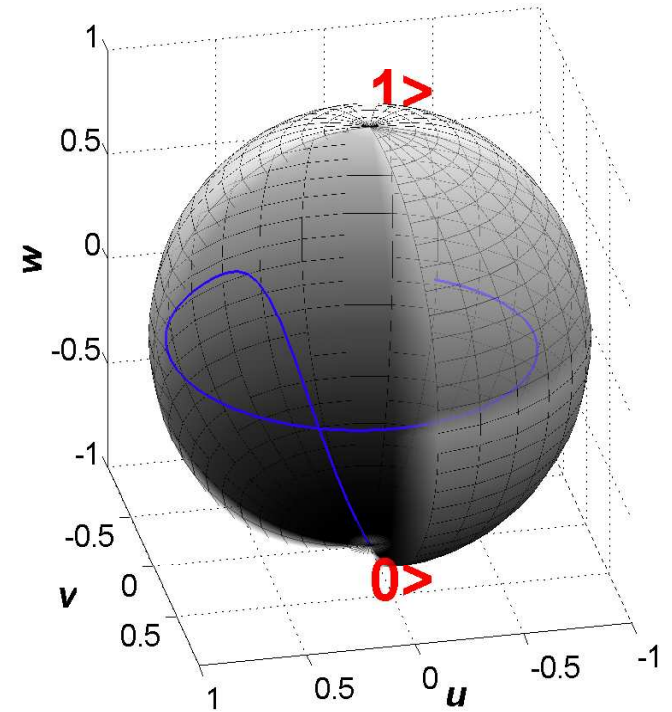
Applications

- ▶ Data storage: 10^{16} bits/cm^3
- ▶ Signal processing: 20+GHz BW

Two-level atoms



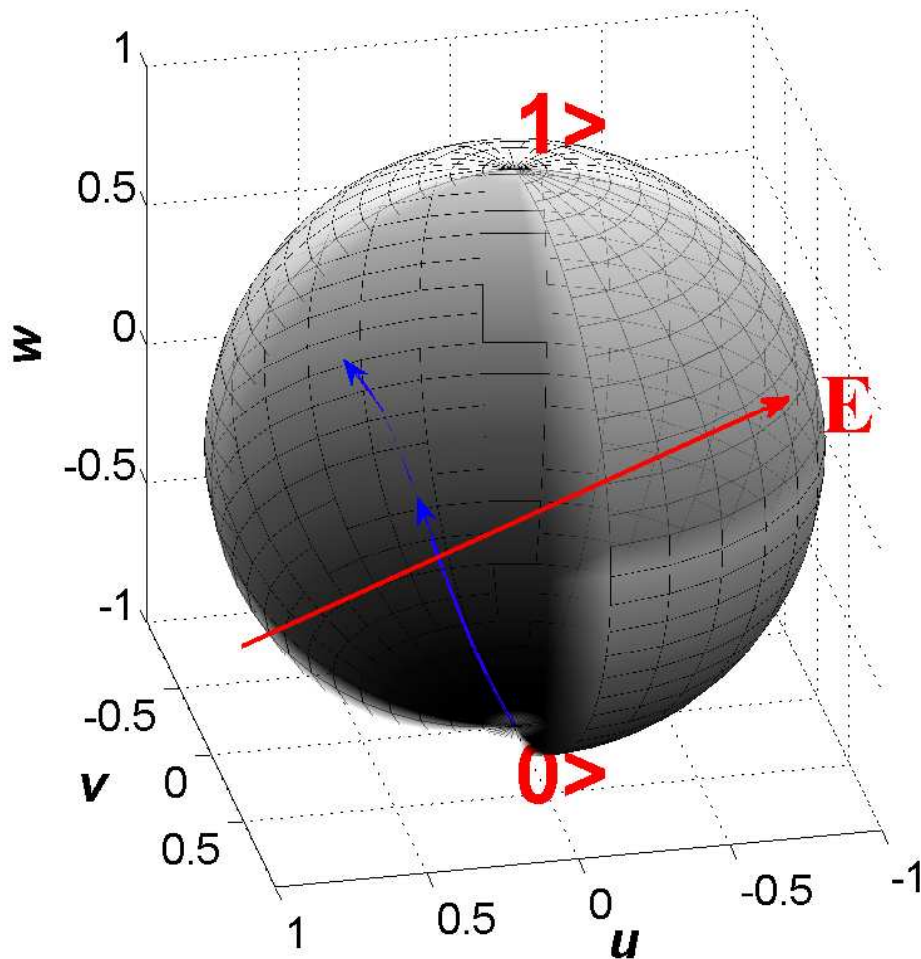
Bloch Sphere



$\text{Tm}^{3+}:\text{YAG}$

Bloch equations: Excitation

Excitation



$$\frac{du}{dt} = -\frac{u}{T_2} - \Delta \cdot v - \kappa \Im\{\mathcal{E}\}w$$

$$\frac{dv}{dt} = \Delta \cdot u - \frac{v}{T_2} + \kappa \Re\{\mathcal{E}\}w$$

$$\frac{dw}{dt} = \kappa \Im\{\mathcal{E}\}u - \kappa \Re\{\mathcal{E}\}v + \frac{w - 1}{T_1}$$

u : inphase component of $\mathcal{P}$

v : in-quadrature component of $\mathcal{P}$

w : population inversion

$\mathcal{E}$: E-field

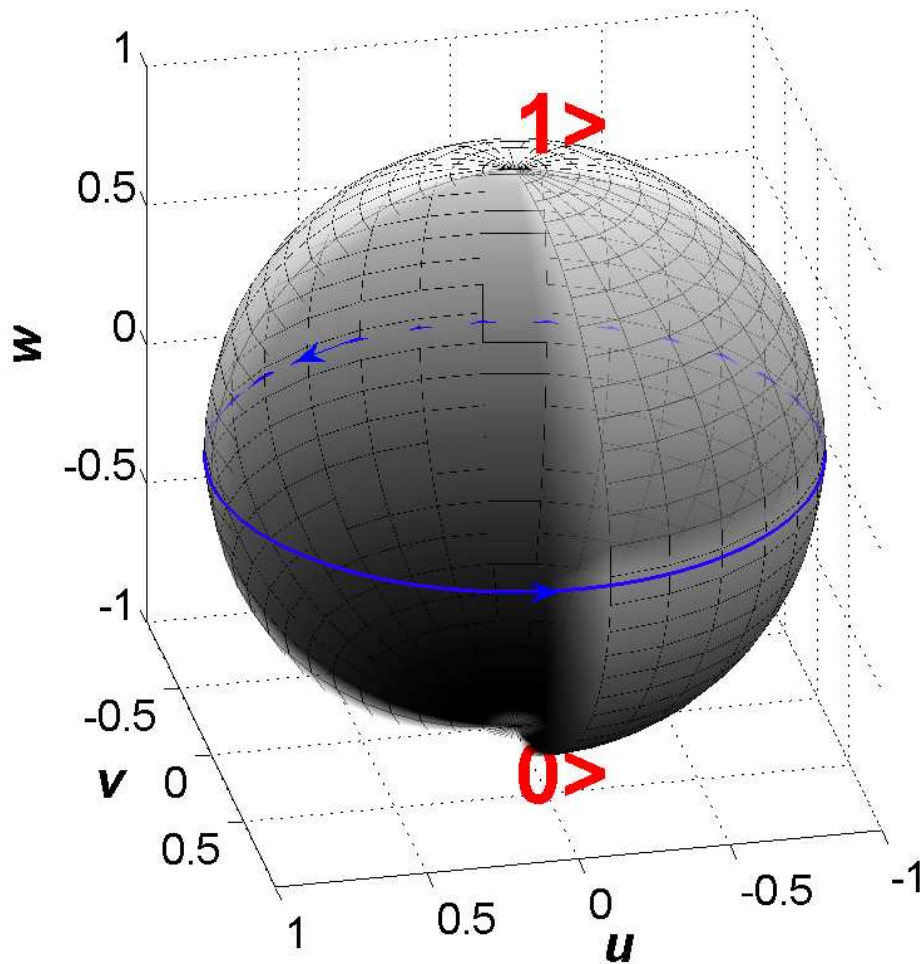
Δ : detuning frequency

T_2 : decoherence time

T_1 : population lifetime

Bloch equations: Detuning

Free precession



$$\frac{du}{dt} = -\frac{u}{T_2} - \Delta \cdot v - \kappa \Im\{\mathcal{E}\}w$$

$$\frac{dv}{dt} = \Delta \cdot u - \frac{v}{T_2} + \kappa \Re\{\mathcal{E}\}w$$

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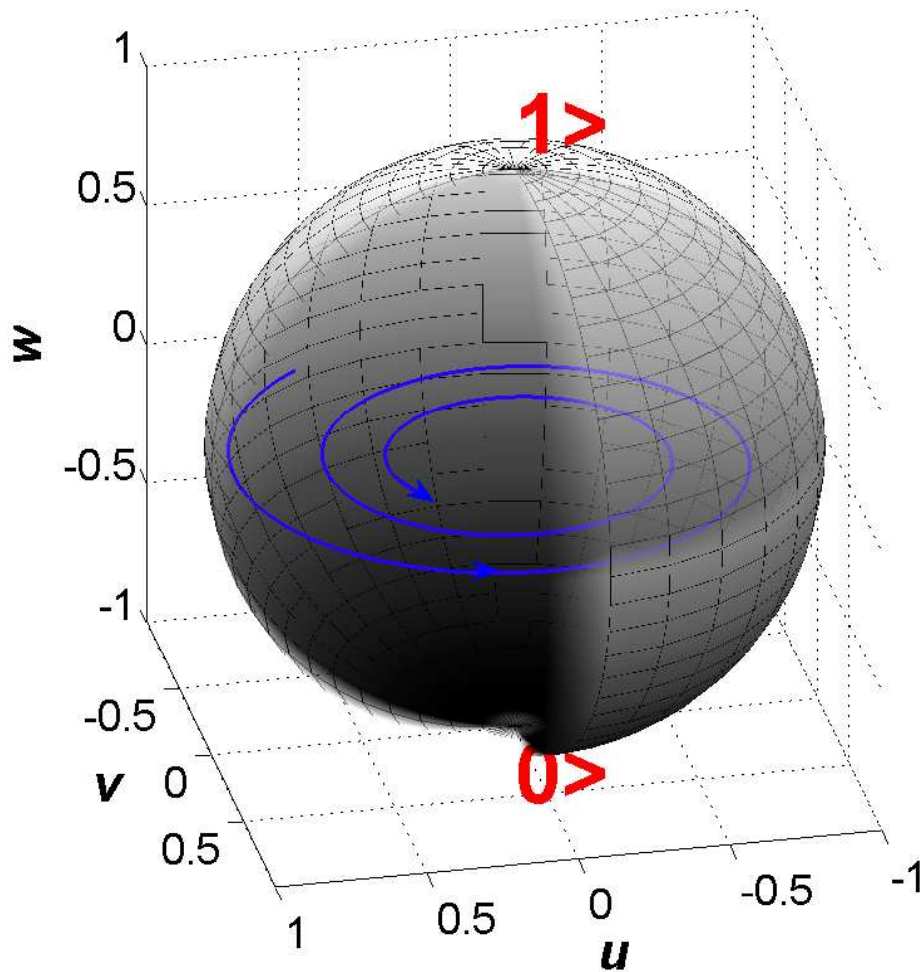
Δ : detuning frequency

T_2 : decoherence time

T_1 : population lifetime

Bloch equations: Decoherence

Decoherence



$$\frac{du}{dt} = -\frac{u}{T_2} - \Delta \cdot v - \kappa \Im\{\mathcal{E}\}w$$

$$\frac{dv}{dt} = \Delta \cdot u - \frac{v}{T_2} + \kappa \Re\{\mathcal{E}\}w$$

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u : inphase component of $\mathcal{P}$

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$\mathcal{E}$: E-field

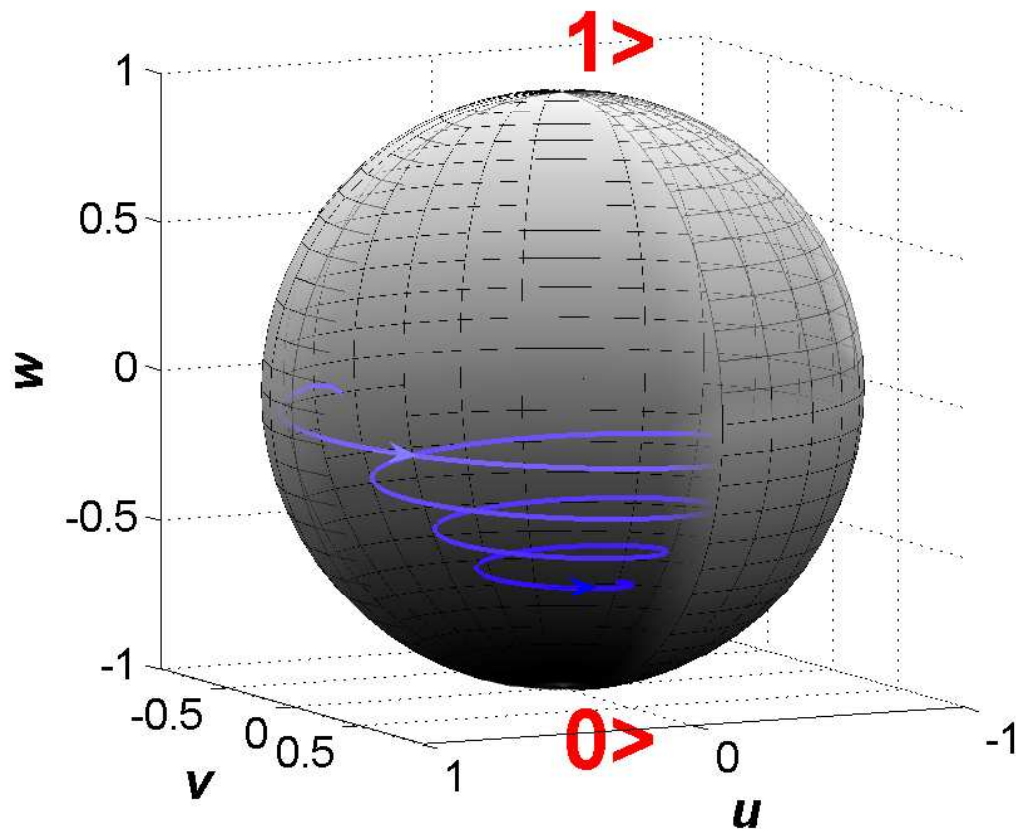
Δ : detuning frequency

T_2 : decoherence time

T_1 : population lifetime

Bloch equations: Population relaxation

Population relaxation



$$\frac{du}{dt} = -\frac{u}{T_2} - \Delta \cdot v - \kappa \Im\{\mathcal{E}\}w$$

$$\frac{dv}{dt} = \Delta \cdot u - \frac{v}{T_2} + \kappa \Re\{\mathcal{E}\}w$$

$$\frac{dw}{dt} = \kappa \Im\{\mathcal{E}\}u - \kappa \Re\{\mathcal{E}\}v + \frac{w - 1}{T_1}$$

u : inphase component of $\mathcal{P}$

v : in-quadrature component of $\mathcal{P}$

w : population inversion

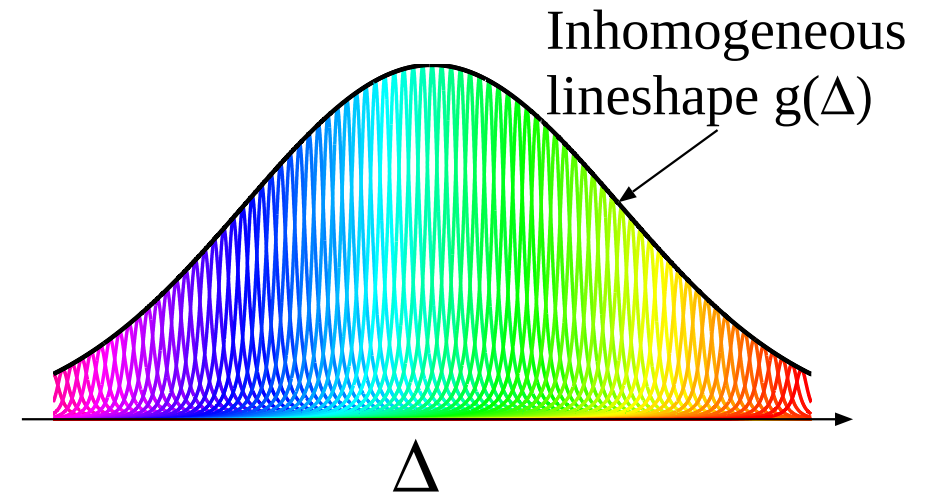
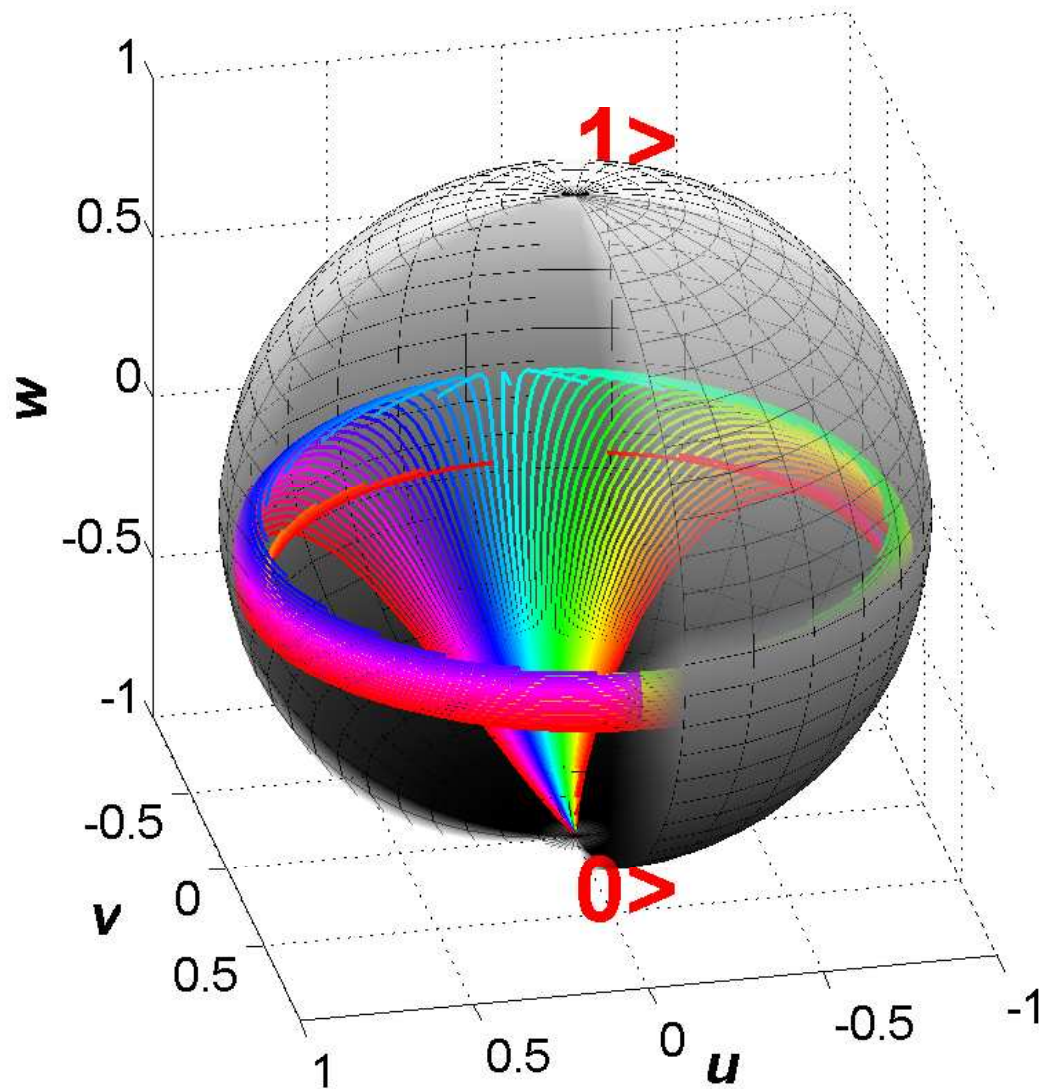
$\mathcal{E}$: E-field

Δ : detuning frequency

T_2 : decoherence time

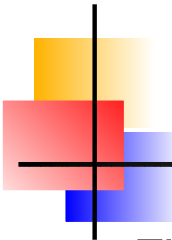
T_1 : population lifetime

Inhomogeneous band



Inhomogeneous averaging

$$\mathcal{P}(t) = \alpha \int_{-\infty}^{+\infty} g(\Delta) [u(t, \Delta) + iv(t, \Delta)] d\Delta$$



2D Maxwell-Bloch equations

The 2D SEVA wave equation is driven by polarization,

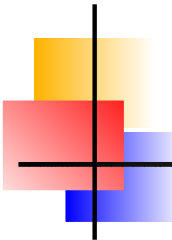
$$\frac{\partial \mathcal{E}(x, z, t)}{\partial z} + \frac{i}{2k} \frac{\partial^2 \mathcal{E}(x, z, t)}{\partial x^2} = -i \frac{\mu_0 \omega_0^2}{2k} \mathcal{P}(x, z, t)$$

Bloch equations under RWA driven by the E-field,

$$\frac{\partial}{\partial t} \begin{bmatrix} u \\ v \\ w \end{bmatrix} = \begin{bmatrix} -\frac{1}{T_2} & -\Delta & -\kappa \Im\{\mathcal{E}\} \\ \Delta & -\frac{1}{T_2} & \kappa \Re\{\mathcal{E}\} \\ \kappa \Im\{\mathcal{E}\} & -\kappa \Re\{\mathcal{E}\} & \frac{1}{T_1} \end{bmatrix} \begin{bmatrix} u \\ v \\ w \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ -\frac{1}{T_1} \end{bmatrix}$$

Polarization given by summation over the inhomogeneous band

$$\mathcal{P}(x, z, t) = \alpha \int_{-\infty}^{+\infty} g(\Delta) [u(x, z, t, \Delta) + iv(x, z, t, \Delta)] d\Delta$$



Previous numerical solutions

▶ Vector form: FDTD^{1,2}

- ▶ Slow: $\Delta z = \frac{\lambda}{100} \sim \frac{\lambda}{50}$ corresponding to $< 0.1\text{fsec}$

2D spatial grid: $200\mu\text{m}(x) \times 5\text{mm}(z)$ 10^8 spatial steps

Time span: $2\mu\text{sec}$ $\Rightarrow 2 \cdot 10^{10}$ time steps

Spectral lines: 250 250 lines

Using 2.25GHz computer, it would take about 3×10^4 years.

▶ Scalar form (wave equation)

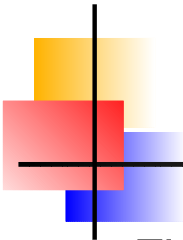
- ▶ Cornish³ simulated 1-D
- ▶ Chang⁴ simulated two symmetric angled plane waves but didn't give specifics of numerical technique

1. Ziolkowski, *et. al*, *Phys. Rev.*, Vol 52, p3082, 1995

2. Schlottau, *et. al*, *Optics Express*, Vol. 13, p182, 2005

3. Cornish, Ph.D. thesis, University of Washington, 2000

4. Chang, *et. al*, *J. of Luminescence*, Vol. 107, p138, 2004



FFT-BPM I

The 2D SEVA wave equation is driven by polarization,

$$\frac{\partial \mathcal{E}(x, z, t)}{\partial z} + \frac{i}{2k} \frac{\partial^2 \mathcal{E}(x, z, t)}{\partial x^2} = -i \frac{\mu_0 \omega_0^2}{2k} \mathcal{P}(x, z, t)$$

Expand $\mathcal{E}(x, z, t)$ and $\mathcal{P}(x, z, t)$ in Fourier domain,

$$\mathcal{E}(x, z) = \int_{-\infty}^{+\infty} \hat{\mathcal{E}}(k_x, z) \exp(-ik_x x) dk_x,$$

$$\mathcal{P}(x, z) = \int_{-\infty}^{+\infty} \hat{\mathcal{P}}(k_x, z) \exp(-ik_x x) dk_x$$

$$-k_x^2 \hat{\mathcal{E}}(k_x, z) - 2ik \frac{d\hat{\mathcal{E}}(k_x, z)}{dz} = -\mu_0 \omega_0^2 \hat{\mathcal{P}}(k_x, z)$$

SEVA wave equation we want to solve is

$$\frac{\partial \hat{\mathcal{E}}}{\partial z} = \frac{i}{2k} (k_x^2 \hat{\mathcal{E}} - \mu_0 \omega_0^2 \hat{\mathcal{P}})$$

Use trapezoidal rule to evolve wave equation

- Since delayed source terms inconsistent with SS-BPM framework

For an ODE $y' = f(x, y)$, the trapezoidal rule is

$$y_{n+1} = y_n + \frac{1}{2}h [f(x_n, y_n) + f(x_n + h, y_{n+1})]$$

Apply trapezoidal rule and solve for $\hat{\mathcal{E}}_{n+1}$,

$$\hat{\mathcal{E}}_{n+1} = \frac{1}{1 - \frac{i}{4k} k_x^2 dz} \left[\left(1 + \frac{i}{4k} k_x^2 dz \right) \hat{\mathcal{E}}_n - \frac{i}{4k} dz \cdot 2\mu_0 \omega_0^2 \hat{\mathcal{P}}_n \right]$$

Numerical scheme for Bloch equations

Bloch equations under RWA driven by the E-field,

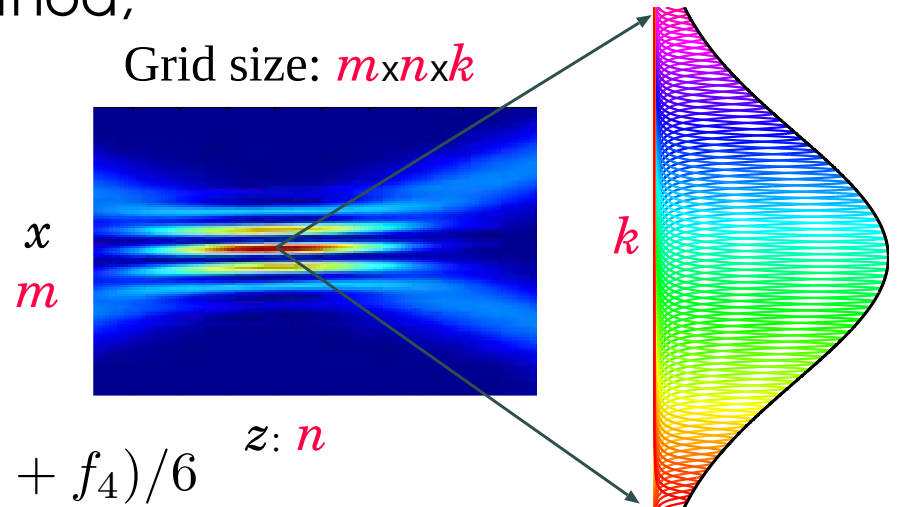
$$\frac{d}{dt} \begin{bmatrix} u \\ v \\ w \end{bmatrix} = \begin{bmatrix} -\frac{1}{T_2} & -\Delta & -\kappa \Im\{\mathcal{E}\} \\ \Delta & -\frac{1}{T_2} & \kappa \Re\{\mathcal{E}\} \\ \kappa \Im\{\mathcal{E}\} & -\kappa \Re\{\mathcal{E}\} & \frac{1}{T_1} \end{bmatrix} \begin{bmatrix} u \\ v \\ w \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ -\frac{1}{T_1} \end{bmatrix}$$

Write in matrix-vector form, it becomes $\underline{y}' = \underline{A}\underline{y} + \underline{b}$

Applying the 4th order Runge-Kutta method,

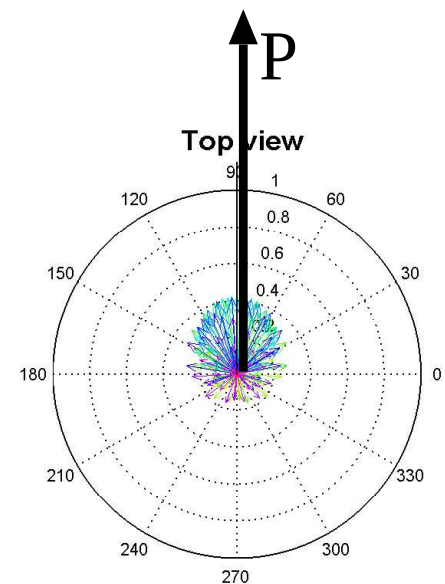
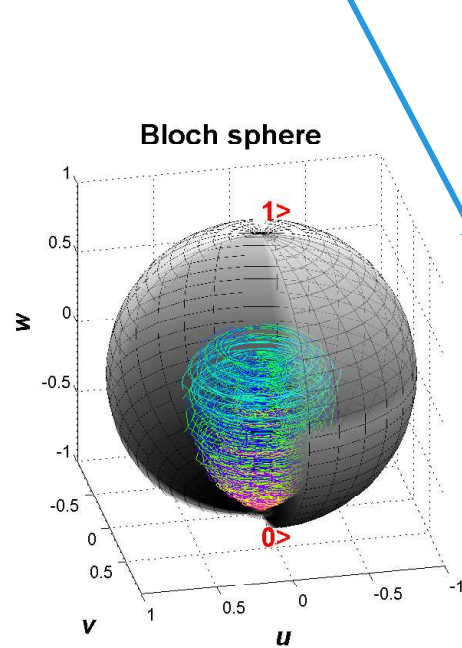
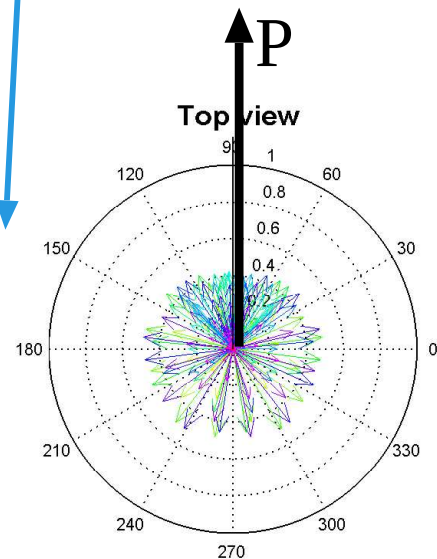
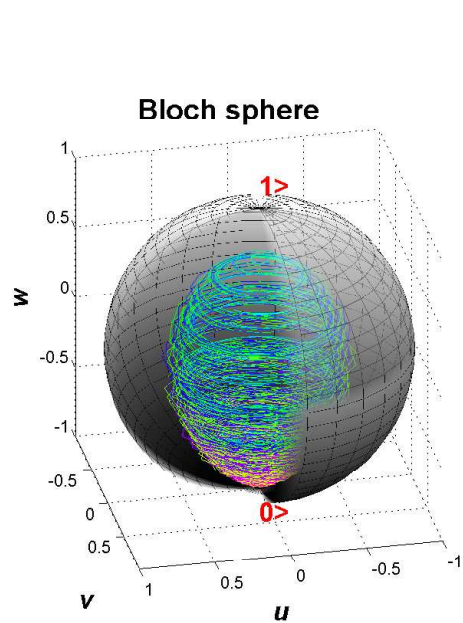
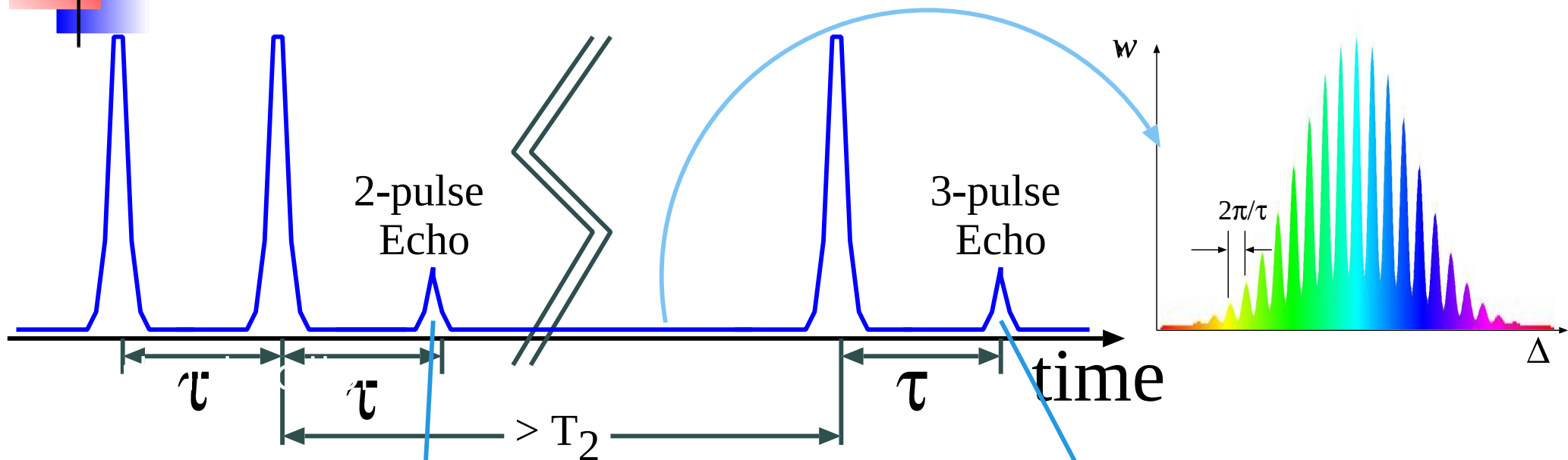
$$\begin{aligned} f_1 &= \underline{A}\underline{y}_{l,m,n,k} + \underline{b}, \\ f_2 &= \underline{A}(\underline{y}_{l,m,n,k} + \frac{f_1}{2}) + \underline{b} \\ f_3 &= \underline{A}(\underline{y}_{l,m,n,k} + \frac{f_2}{2}) + \underline{b}, \\ f_4 &= \underline{A}(\underline{y}_{l,m,n,k} + f_3) + \underline{b} \end{aligned}$$

$$\underline{y}_{l+1,m,n,k} = \underline{y}_{l,m,n,k} + \Delta t(f_1 + 2f_2 + 2f_3 + f_4)/6$$



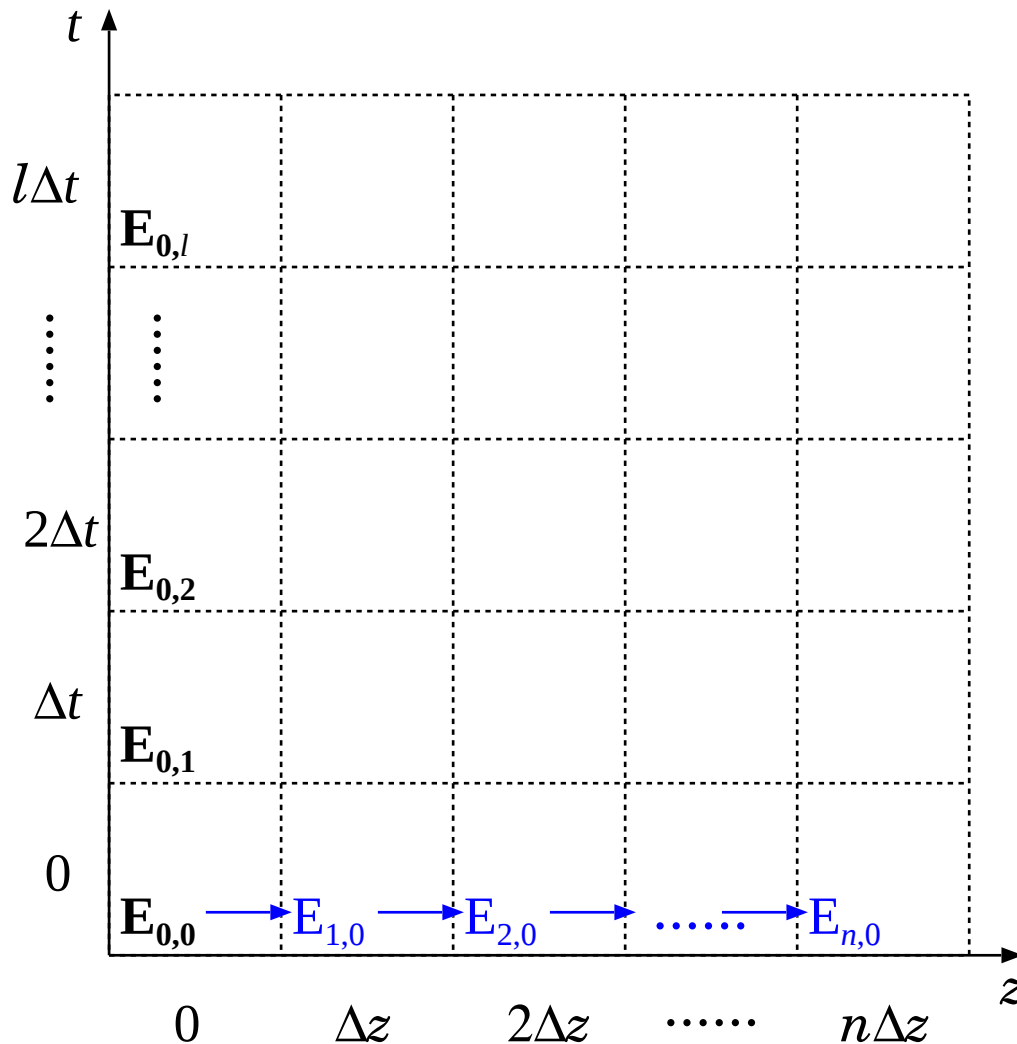
l : time step, m : x sampling step, n : z step, k : spectral lines

Bloch simulation results

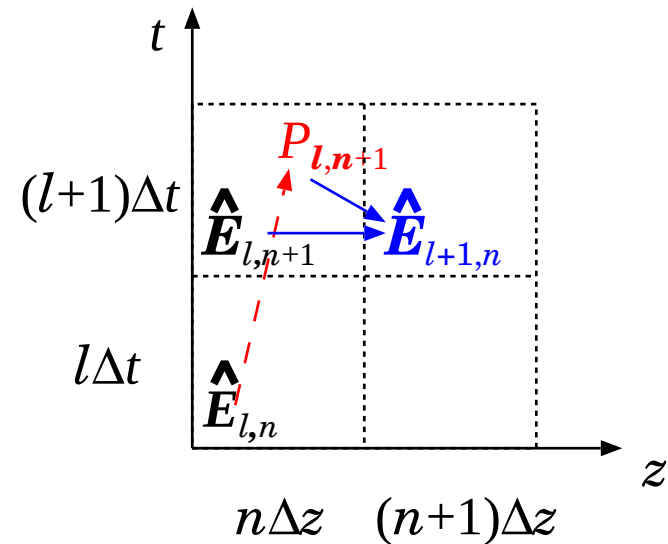


Advancing the Maxwell-Bloch equations

Initial Value Condition

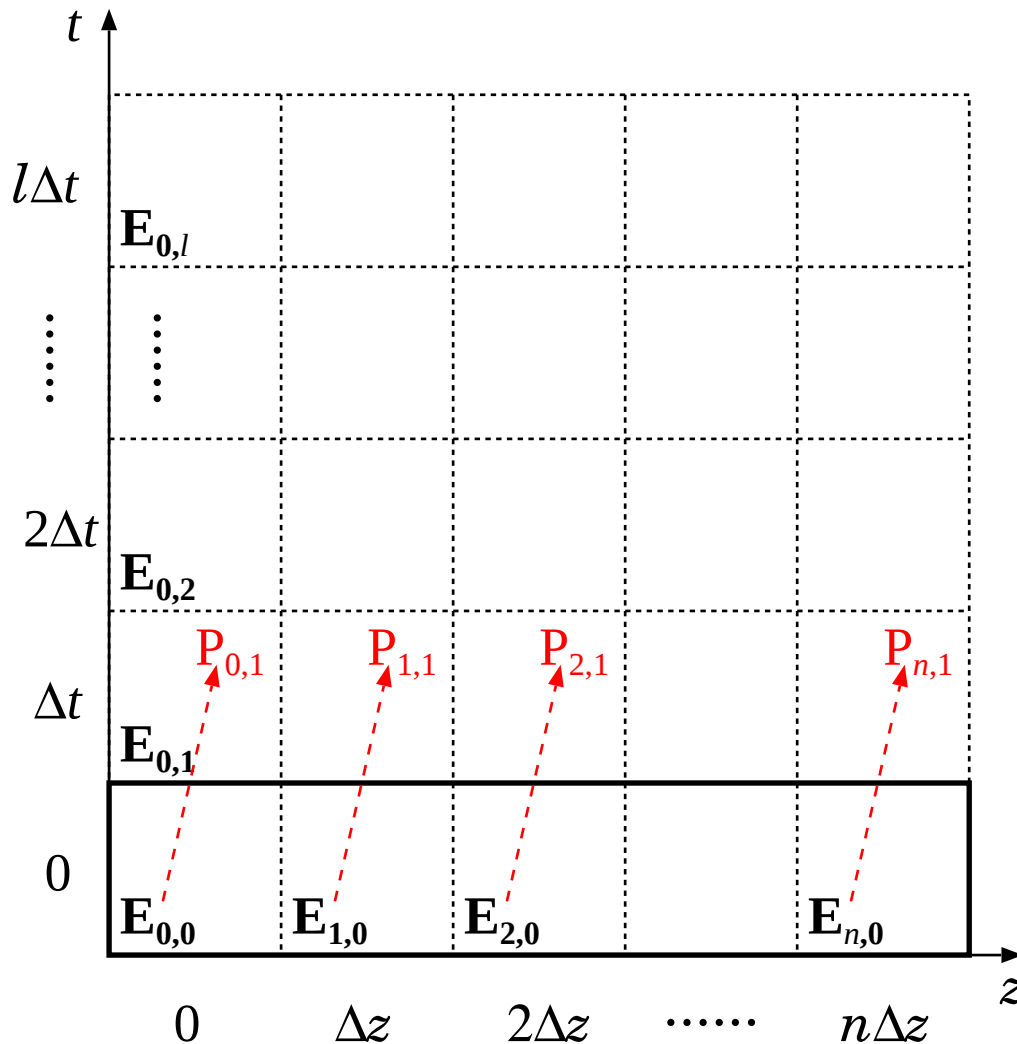


- ▶ Calculate **Maxwell's** equation
- ▶ Propagate E-field from boundary into crystal

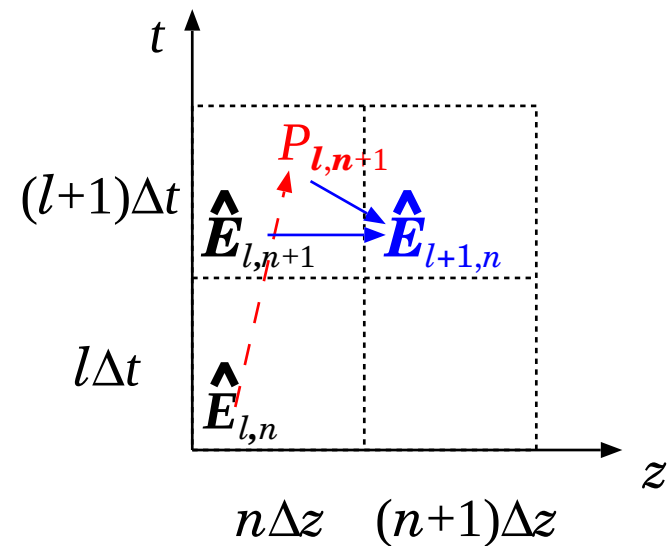


Advancing the Maxwell-Bloch equations

Δt later

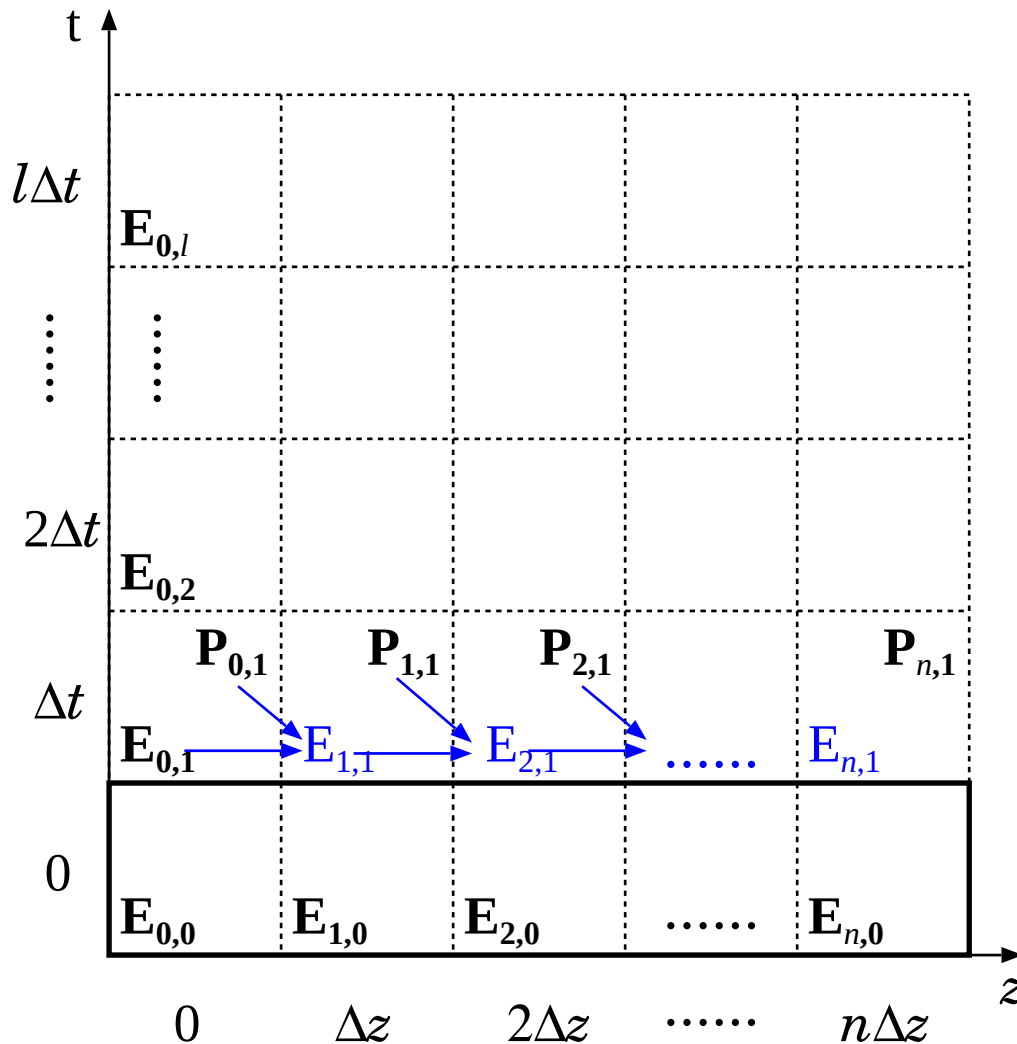


- Calculate **Bloch** equation
- E-field drives Bloch vectors to get polarization

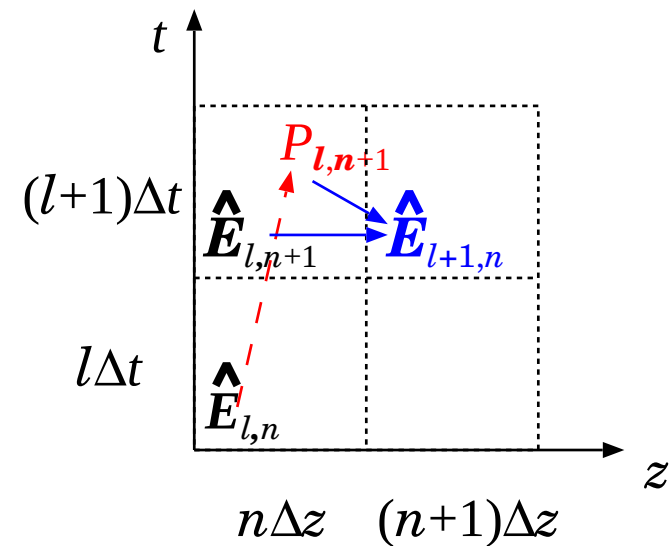


Advancing the Maxwell-Bloch equations

Δt later

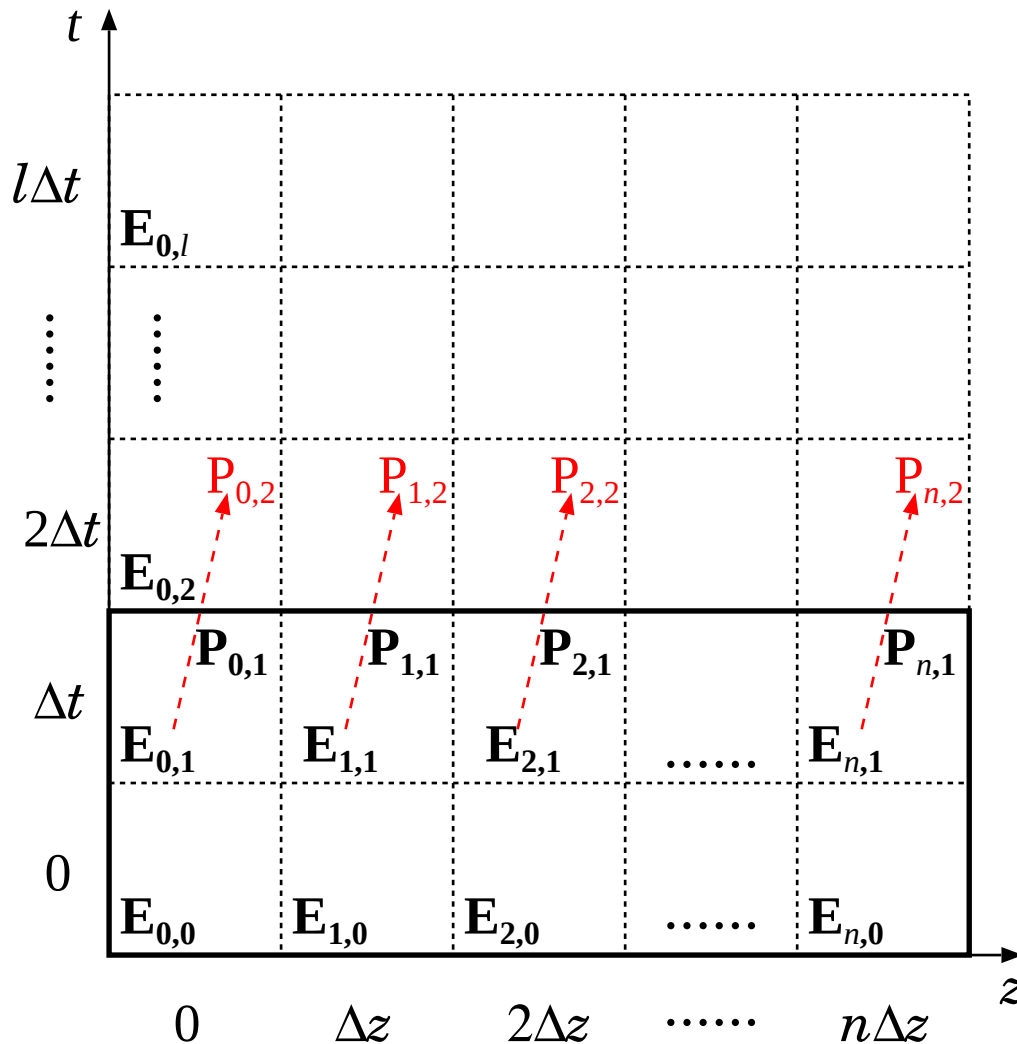


- ▶ Calculate **Maxwell's** equation
- ▶ Contribute polarization to propagate E-field

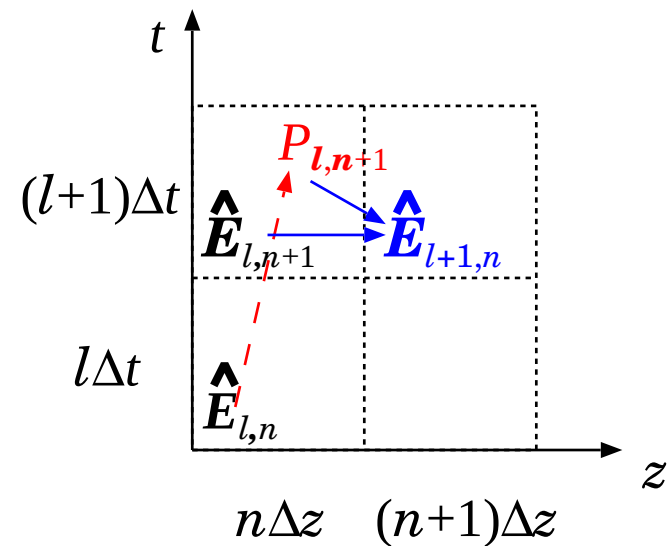


Advancing the Maxwell-Bloch equations

$2\Delta t$ later

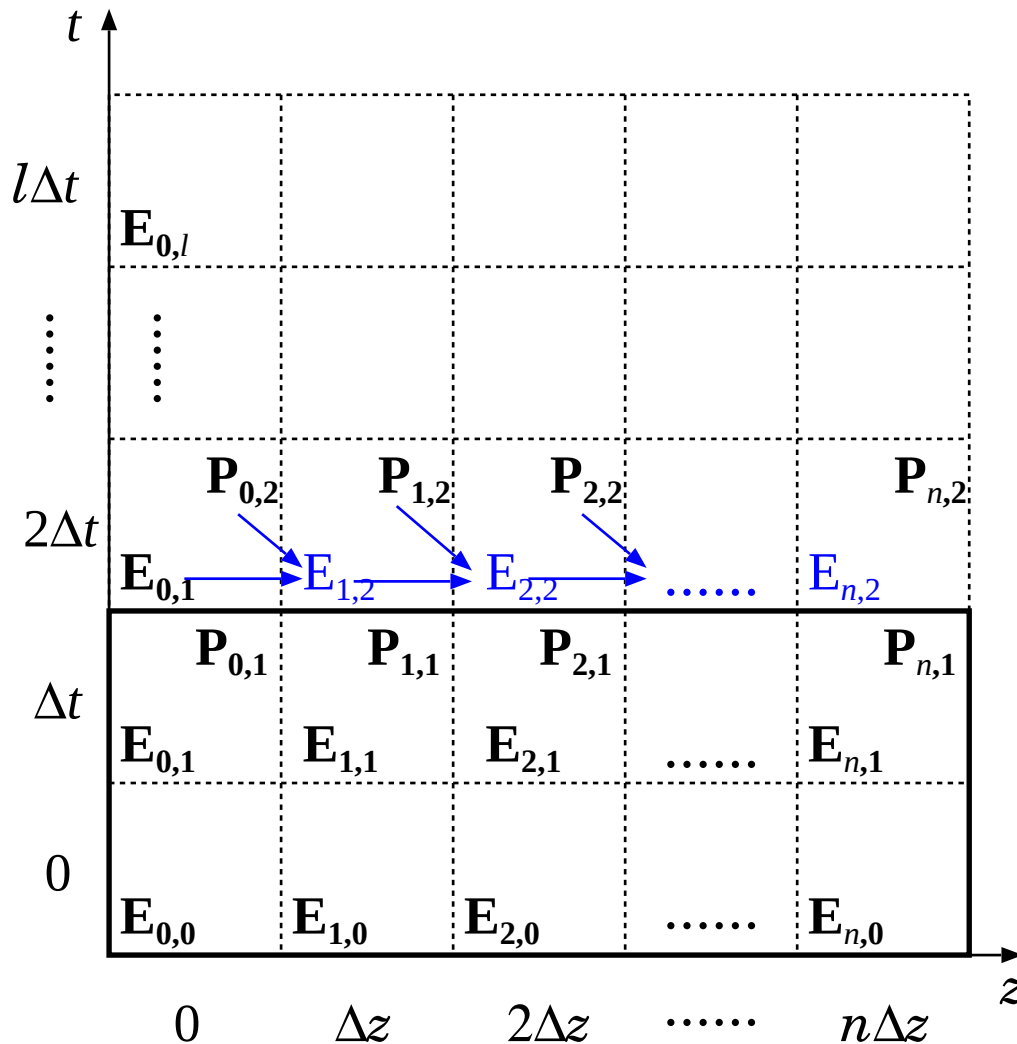


- Calculate **Bloch** equation
- E-field drives Bloch vectors to get polarization

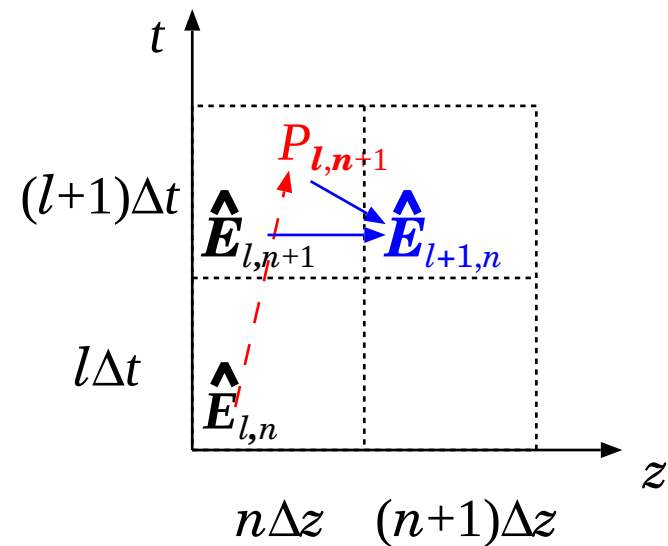


Advancing the Maxwell-Bloch equations

$2\Delta t$ later

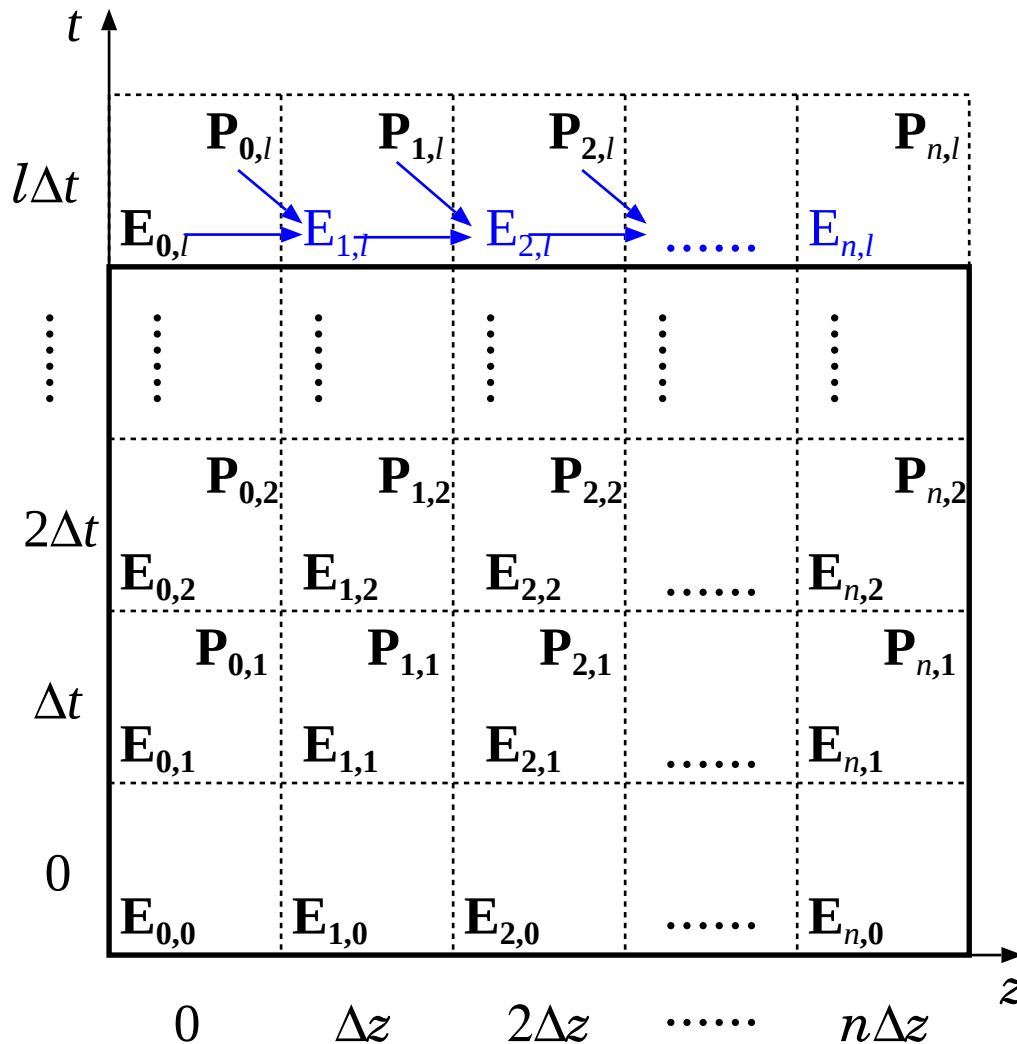


- ▶ Calculate **Maxwell's** equation
- ▶ Contribute polarization to propagate E-field

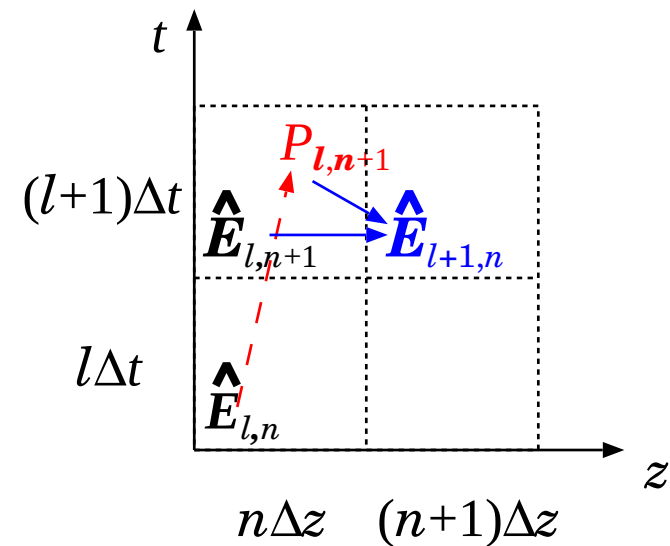


Advancing the Maxwell-Bloch equations

$m\Delta t$ later

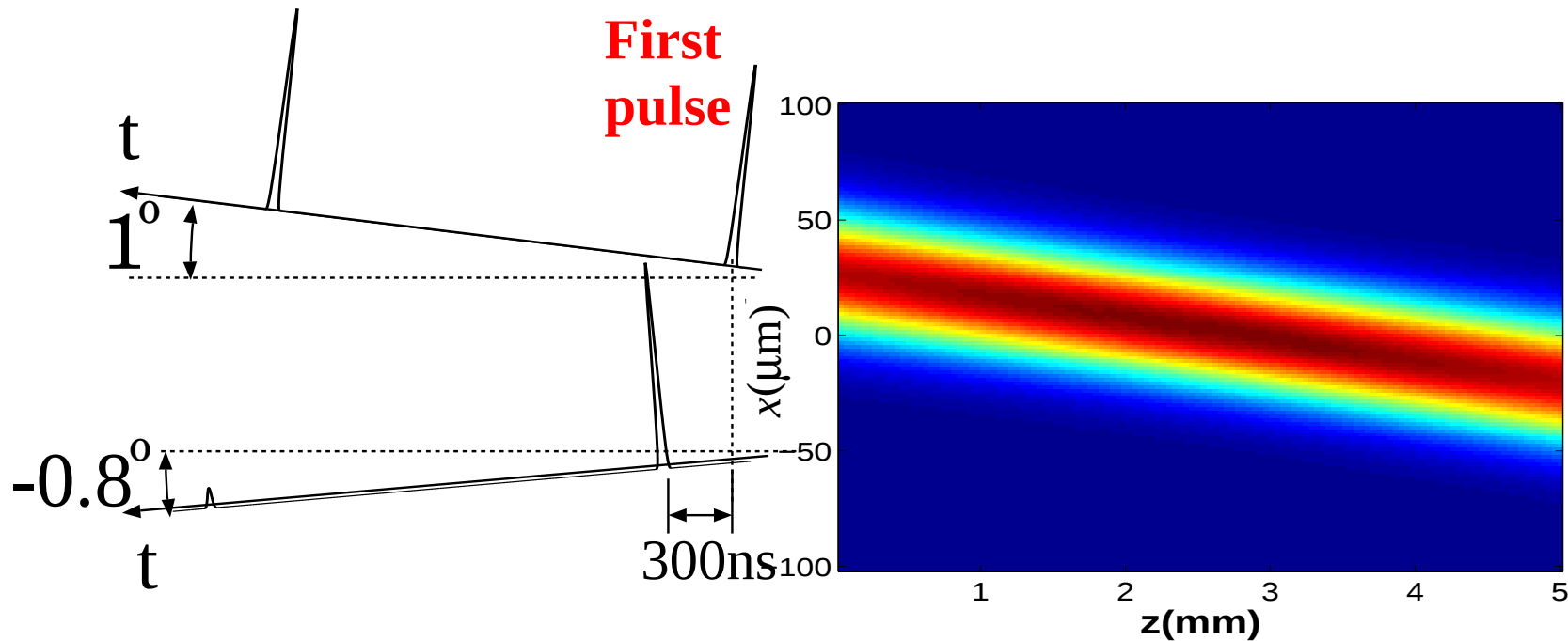


- Calculate **Maxwell's** equation
- Contribute polarization to propagate E-field



Maxwell-Bloch simulation setup

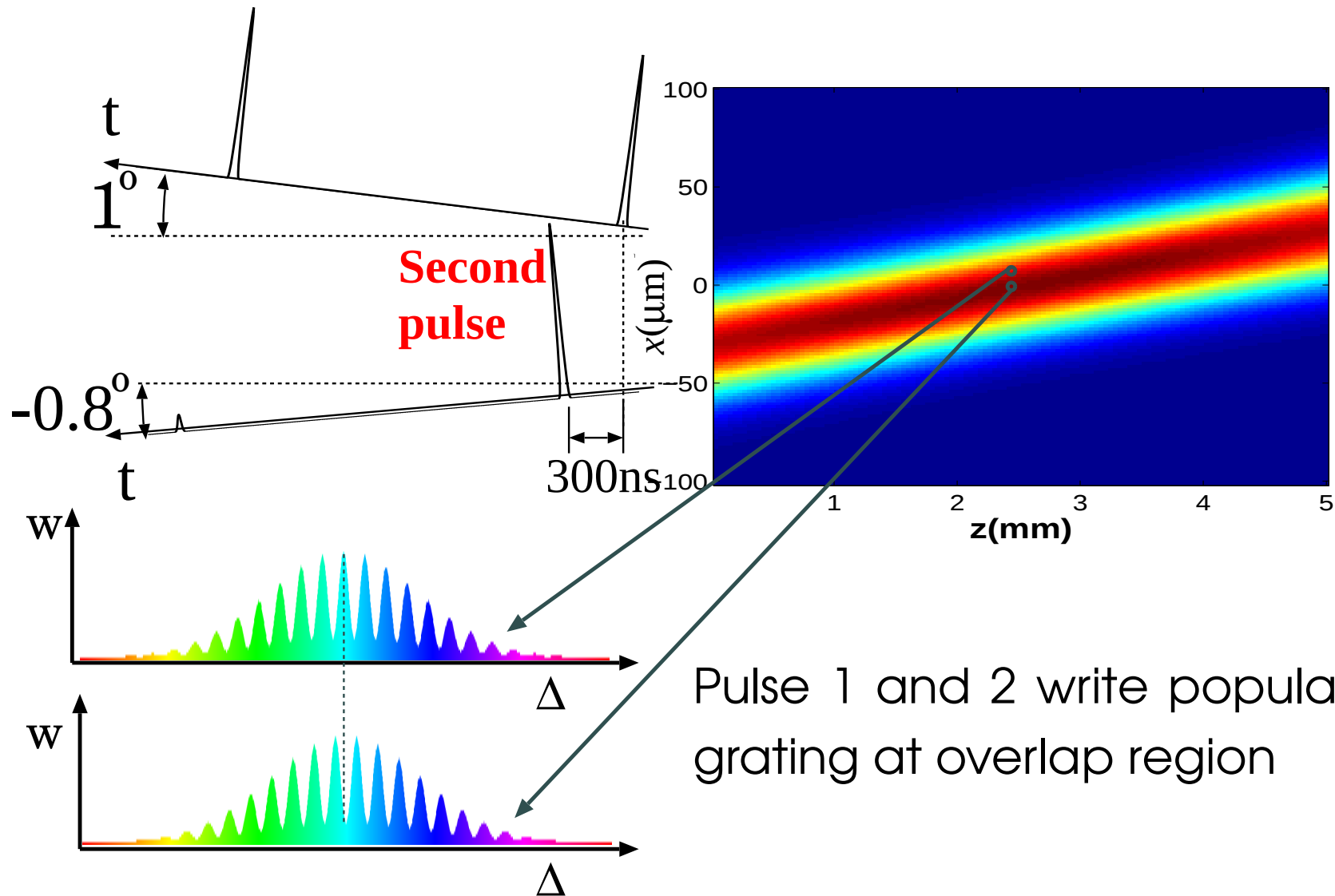
$200\mu\text{m} \times 5\text{mm}$, $2.4\mu\text{sec}$, 250 spectral lines, 100MHz of BW



Pulse 1 comes in at 1°

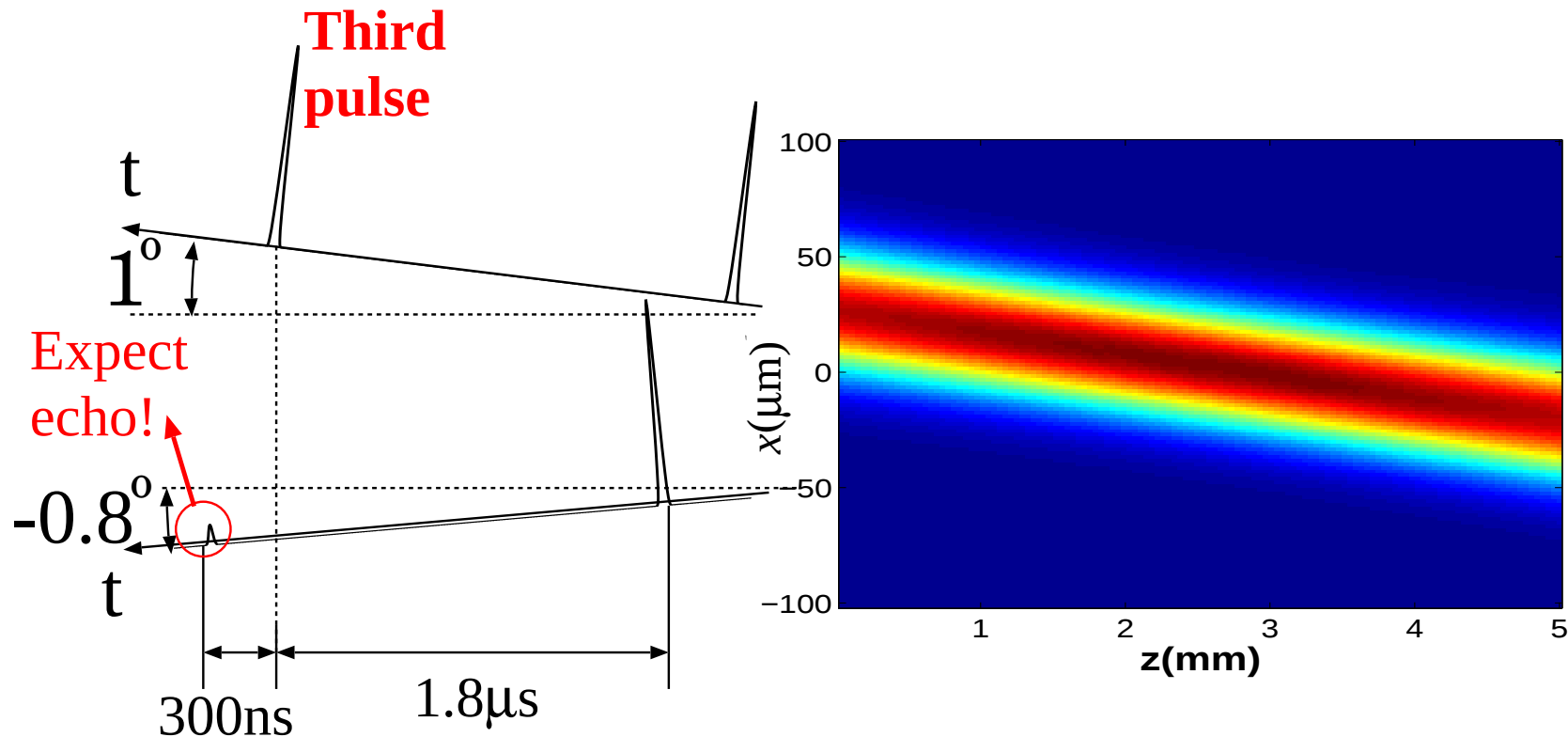
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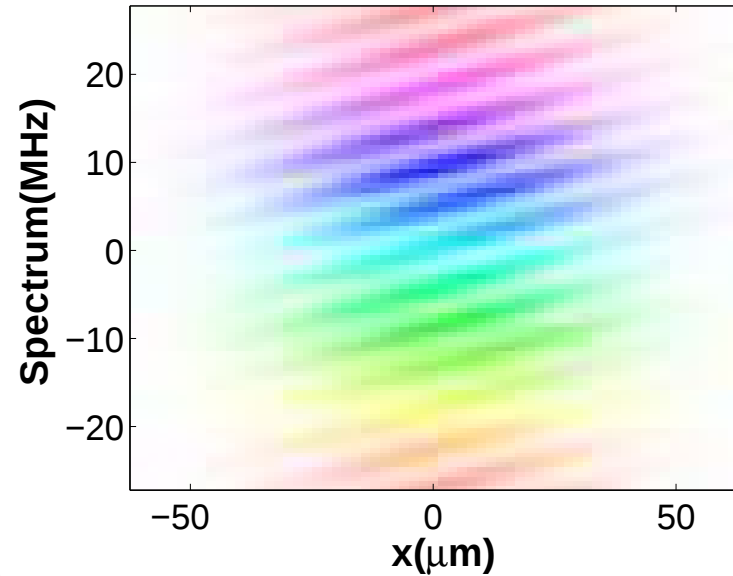
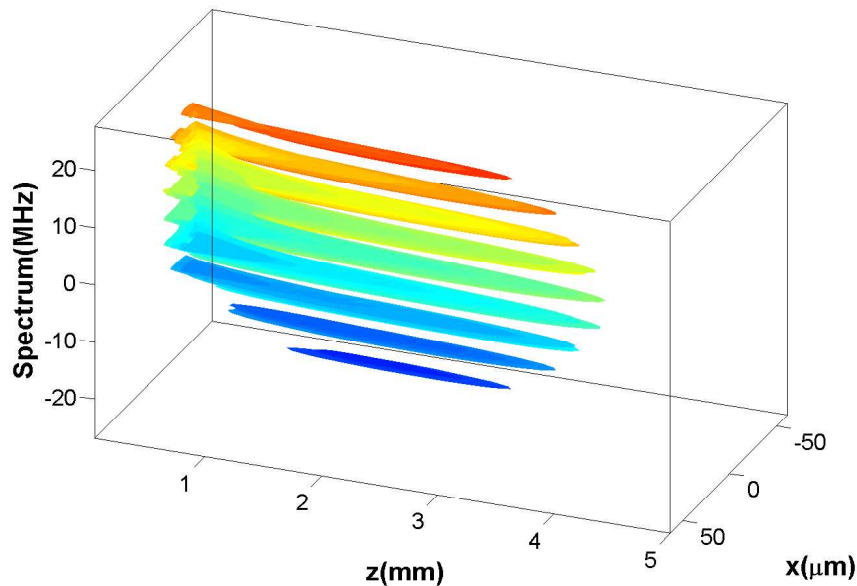
$200\mu\text{m} \times 5\text{mm}$, $2.4\mu\text{sec}$, 250 spectral lines, 100MHz of BW



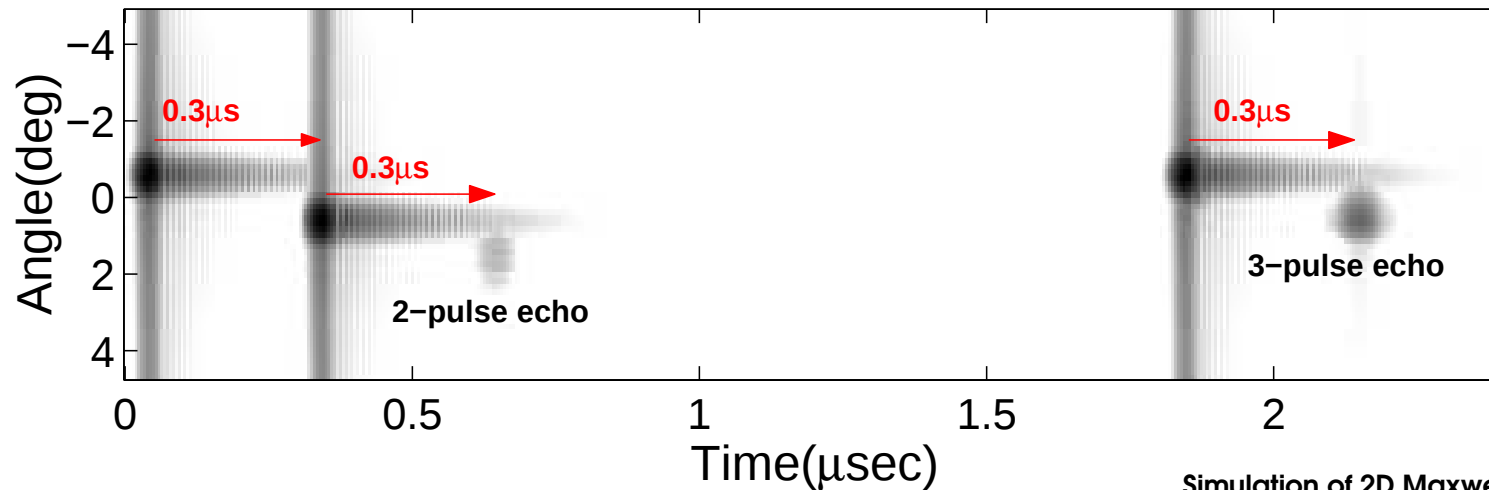
- ▶ Pulse 3 comes in at pulse 1 direction
- ▶ Population gating diffracts pulse 3 both in space and time
- ▶ Echo is expected 300ns later in direction of pulse 2

Angled beam photon echo

Volume population grating

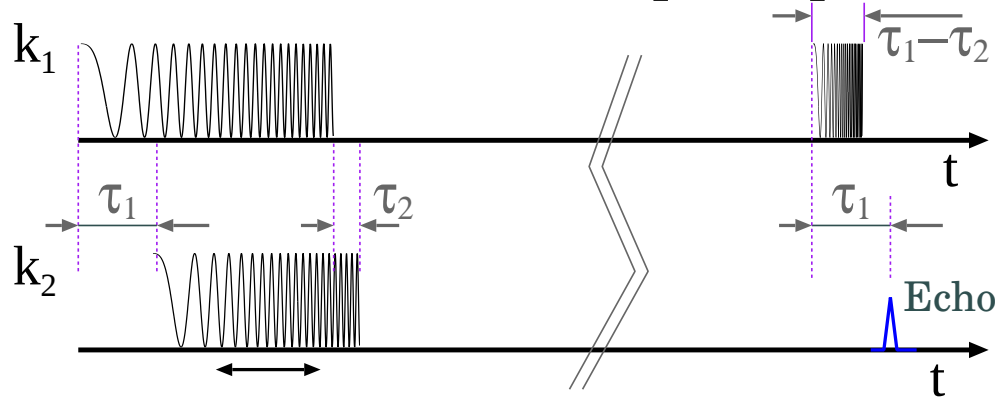


Time-angle waterfall display

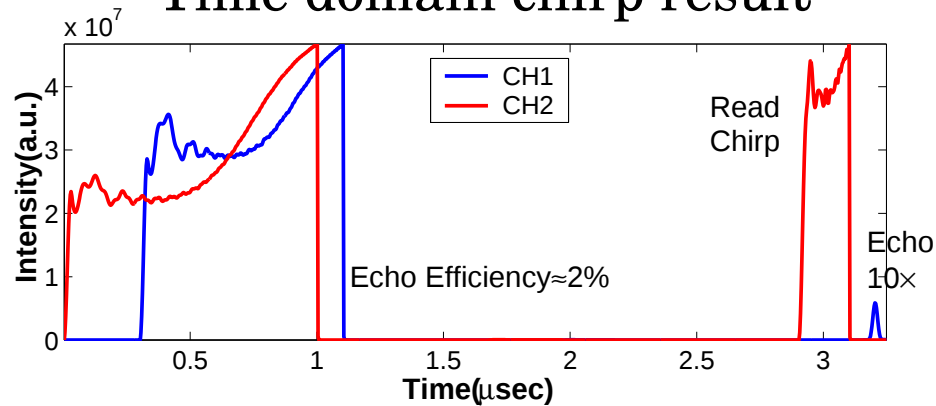


Chirped angled beam photon echoes

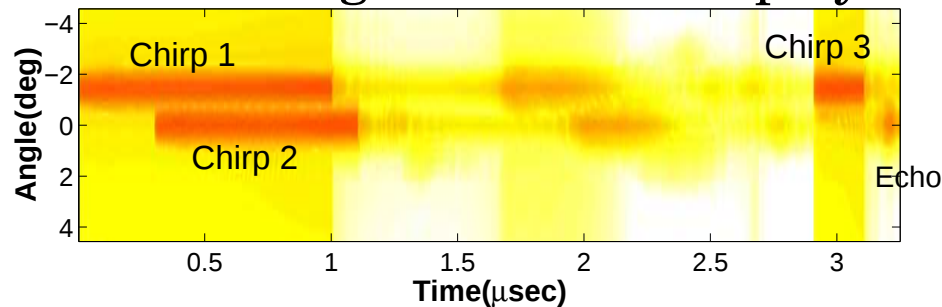
Time domain chirp setup



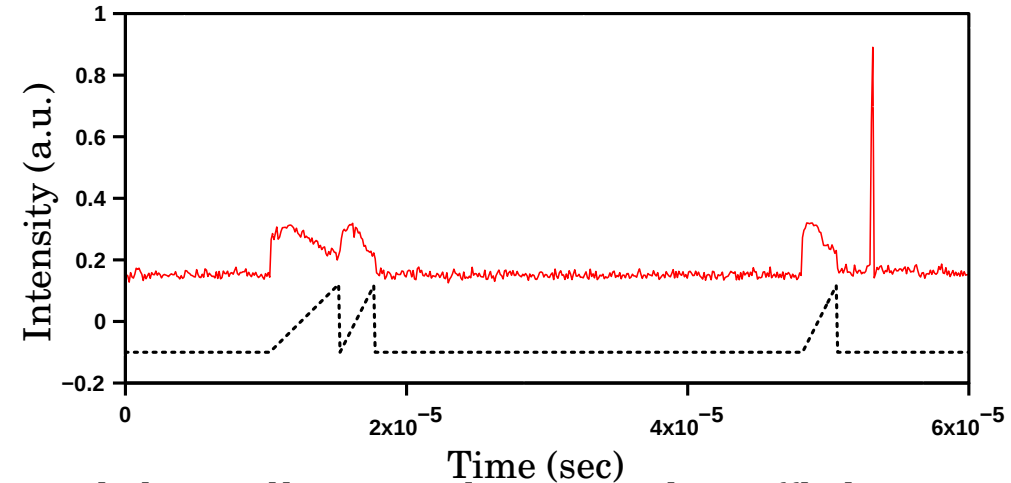
Time domain chirp result



Time-angle waterfall display

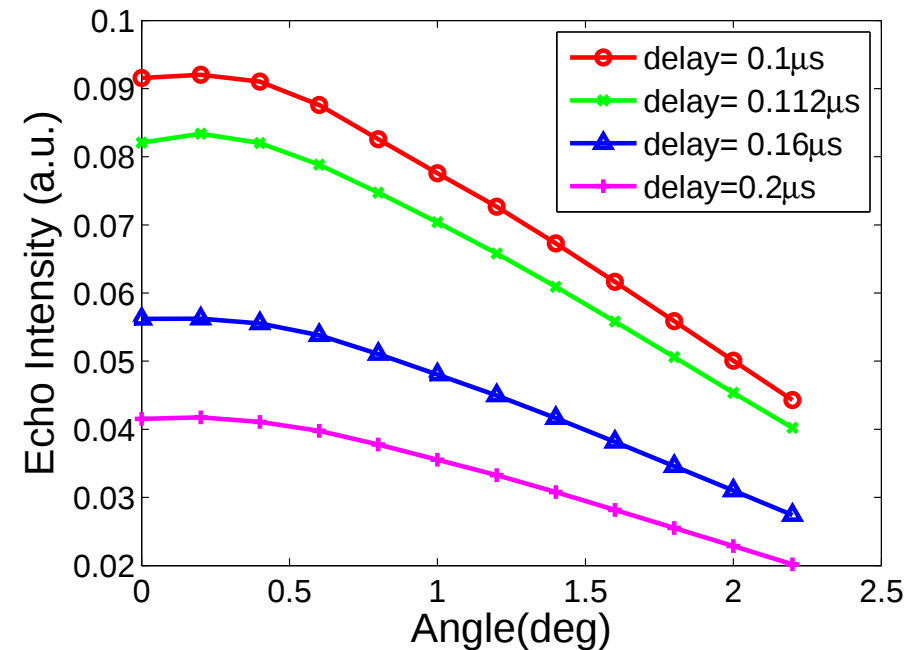


Chirp photon echo experimental traces



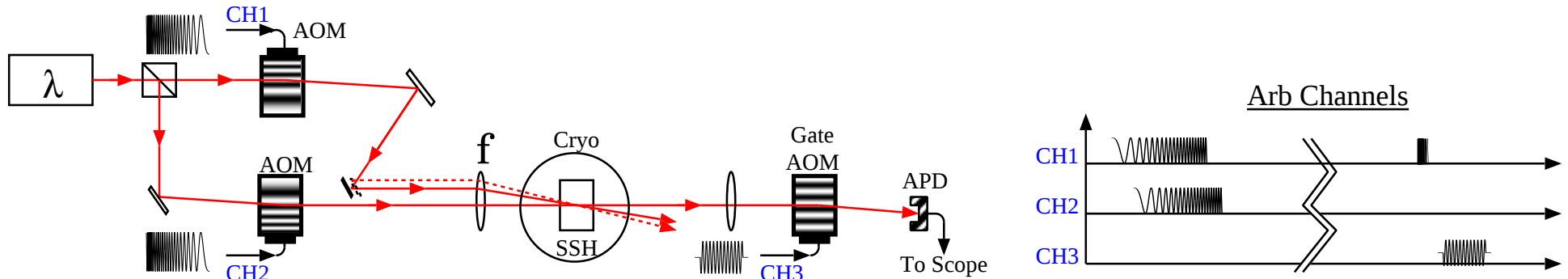
Interaction angle vs. echo efficiency

$w_0 = 60 \mu\text{m}$, $\alpha L = 1.0$



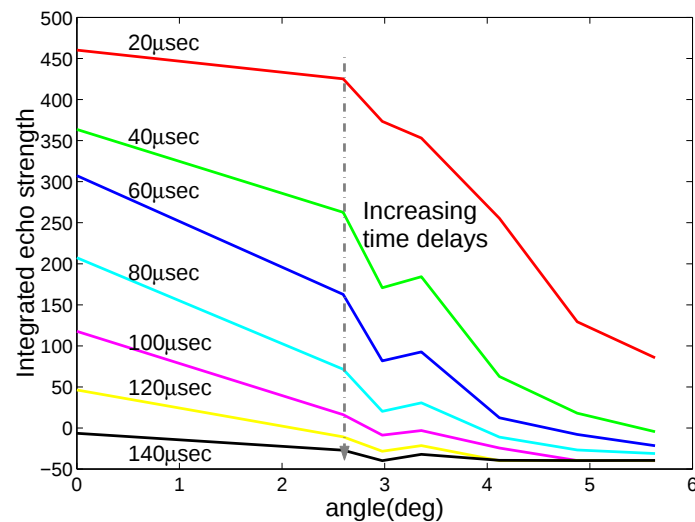
Experimental results

Experimental Setup



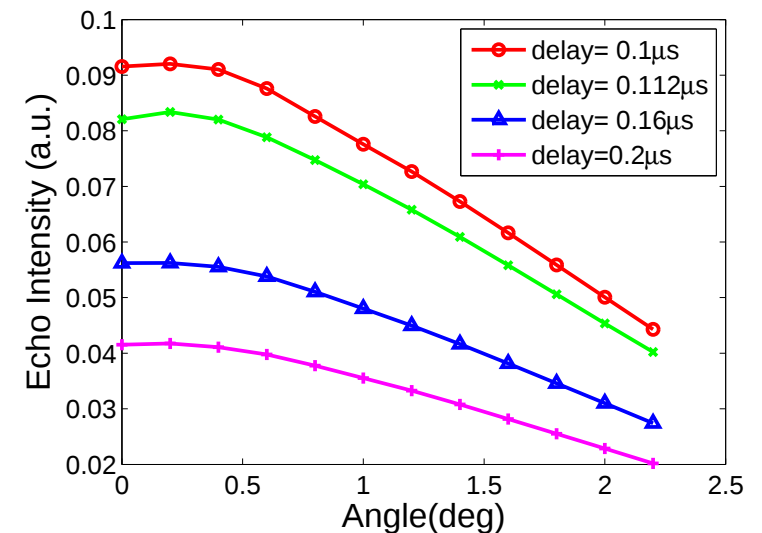
Experimental Results

Spot size: $40\mu\text{m} \times 80\mu\text{m}$



Simulation Results

$w_0 = 60\mu\text{m}$, $\alpha L = 1.0$



Conclusion

- ▶ Spatial-spectral Maxwell-Bloch equations solved using FFT-trapezoidal rule and RK-4

- ▶ We are now working towards
 - ▶ Increase the number of transverse samples for larger angles
 - ▶ Simulate many GHz of bandwidth
 - ▶ Simulate realistic T_2 (T_1 effect for accumulation)

Acknowledgement

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