IEEE NUSOD 2007, Newark, DE



Simulation of Derivative Characteristics of Broadband Quantum Dot Lasers

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Acknowledgement: National Science Foundation (Grant No.: ECCS-0725647)







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Introduction



Applications of Broad Gain Material & Broadband Laser

Optical Telecommunications

- Ultra-broadband components
 - tunable laser, SOA, EA modulator, detector, etc
- Ultra-short pulse generation
 - Optical clocking , OTDM, etc

Spectroscopy & Sensing

- Molecular spectroscopy (1450-1650nm)
 - Strong overtone spectra of CO, C_2H_2 , and NH_3
- Atmospheric and planetary gas sensors
 - CH_4 , CO, CO_2 , H_2S , HCI, NH_3 , C_2H_4 , C_2H_2 , C_2H_6 , C_6H_6 , etc
- General Spectroscopy
 - Material absorption, transmission, luminescence, etc

Metrology

- Optical test and measurements, etc
- Optical time domain reflectrometry (OTDR)

Imaging

- Bio-imaging (Optical Coherence Tomography)
- Ultra-short pulse imaging, etc

Others

- High sensitive fiber gyroscope
- Instrumentation, etc





LEHIGH UNIVERSITY Introduction (con't) Center for Optical Technologies

NATURE VOL 415 21 FEBRUARY 2002 www.nature.com

Ultra-broadband semiconductor laser

Claire Gmachl, Deborah L. Sivco, Raffaele Colombelli, Federico Capasso & Alfred Y. Cho

Bell Laboratories, Lucent Technologies, 600 Mountain Avenue, Murray Hill, New Jersey 07974, USA





Intersub-band cascade → mid-IR
Quantum band engineering:

36 different active regions

Covering 6-8 µm emission
Low wall-plug efficiency at RT (<0.1%)
Side-mode-supression-ratio : ~20 dB
Material challenge for <u>near-IR region</u>!











Theoretical Model

$$\frac{dN_{w}}{dt} = \frac{\eta_{i}I}{q} - \frac{N_{w}}{T_{wr}} - \frac{N_{w}}{T_{wu}}$$

$$- \frac{N_{e}}{T_{e}} + \frac{\sum N_{u,j}}{T_{uw}}$$

$$\frac{dN_{u,j}}{dt} = \frac{N_{w}G_{n}}{T_{wu,j}} + \frac{N_{g,j}}{T_{gu,j}} + \frac{N_{e,j}}{T_{eu,j}} - \frac{N_{u,j}}{T_{ug,j}} - \frac{N_{u,j}}{T_{ue,j}}$$

$$- \frac{N_{u,j}}{T_{uw}} - \frac{N_{u,j}}{T_{r}} - \frac{N_{u,j}}{T_{e}} - \frac{C\Gamma}{n_{r}} \sum_{m} g_{mn} Sm$$

$$\frac{dN_{e,j}}{dt} = \frac{N_{u,j}}{T_{r}} - \frac{N_{e,j}}{T_{e}} - \frac{C\Gamma}{n_{r}} \sum_{m} g_{mn} S_{m}$$

$$\frac{dN_{g,j}}{dt} = \frac{N_{u,j}}{T_{e}} + \frac{N_{e,j}}{T_{eg,j}} - \frac{N_{g,j}}{T_{r}} - \frac{N_{e,j}}{T_{gu,j}} - \frac{N_{e,j}}{T_{eg,j}} - \frac{N_{g,j}}{T_{gu,j}} - \frac{N_{g,j}}{T_{ge,j}} - \frac{N_{g,j}}{$$



LEHIGH Theoretical Model (con't)



Optical gain modal:

$$g_{mn} = \frac{2 \pi q^{2} \hbar D_{l} N_{D}}{cn_{r} \varepsilon_{0} m_{0}^{2}} \frac{|P_{cv}^{\sigma}|^{2}}{E_{cv,k_{l}}} (2 P_{n} - 1) G_{n} B_{cv} (E_{m} - E_{cv,k_{l}}^{j})$$

$$P_{l,j} = \frac{N_{l,j}}{2 D_{l} N_{D} G_{n}}$$

Homogeneous Broadening function:

$$B_{cv} (E_m - E_{cv,k_l}^{j}) = \frac{\hbar \Gamma_l^{j} / \pi}{(E_m - E_{cv,k_l}^{j})^2 + (\hbar \Gamma_l^{j})^2}$$

Kramers-Kronig relation to Δn in QD ensemble:

$$\Delta n_{QD}(E) = k \frac{\hbar c}{2E} \Gamma N_D \sum_{l} \sum_{j} D_l \frac{|P_l^{\sigma}|^2}{E_l} [(2P_{l,j} - 1)G_n \frac{(E - E_{cv,l})/\pi}{(E - E_{cv,l})^2 + (\hbar \Gamma_{cv,l})^2}]$$

M. Sugawara et al., Phys. Rev. B, 61, 11 (2000).

M. Sugawara et al., J. Appl. Phys., 97, 043523 (2005).

M. Gioannini et al., Opt. and Quantum Electron., 38, 381-394 (2006).

Theoretical Model (con't)



Contribution of free carriers in the wetting layer to Δn :

$$\Delta n_{WL} = \Gamma_{WL} K_n \frac{\Delta N_{WL}}{E^2}$$

Linewidth Enhancement Factor below threshold:

 $\alpha = -\frac{4\pi}{\lambda} \frac{\left(\Delta n_E + \Delta n_G + \Delta n_{UC} + \Delta n_{WL}\right) / \Delta N}{\Delta g / \Delta N}$

Linewidth Enhancement Factor above threshold:

$$\alpha = \alpha_{QD,E} + \alpha_{QD,G} + (\alpha_{UC} + \alpha_{WL}) \frac{\rho_{th} - 1 + \rho_{th} \frac{J}{J_{th}}}{2\rho_{th} - 1}$$

Y. Toda *et al.*, Phys. Rev. Lett., 82, 20 (1999).
M. Gioannini *et al.*, Opt. and Quantum Elec., 38, 381-394 (2006).
S. Melnik *et al.*, Opt. Exp., 14, 2950-2955 (2006).
Z. Mi *et al.*, IEEE J. Quan. Elec., 43, 5 (2007).





Simulation Results



• The calculated alpha of typical QD lasers (Type I) at below and above threshold shows similar trend as reported from literature reviews.

• QD lasers with broadband nature (Type III) shows larger change of alpha above threshold in accordance to injection as compared to typical QD lasers. This is due to the larger alpha contribution from continuum states in large inhomogeneous system, as can be found from following figures.

S. Melnik *et al.*, Opt. Exp., 14, 2950-2955 (2006). Z. Mi *et al.*, IEEE J. Quantum. Electron., 43, 5 (2007).



Simulation Results (con't)



• Broadband QD lasers (Type III) shows a smaller degree of change over energy span as compared to typical QD lasers (Type I) from figure.

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• The arrows show the lasing energy of each system. The up-arrow refers to Type I while down-arrow refers to Type II and III.

• The LEF of typical QD lasers shows value below 2 near threshold.

J. Oksanen *et al.*, J. Appl. Phys., 94, 3 (2003).
M. Gioannini *et al.*, Opt. and Quantum Elec., 38, 381 (2006)
S. Melnik *et al.*, Opt. Exp., 14, 2950-2955 (2006).





• Typical QD lasers show that alpha contribution from continuum states increase rapidly with injection as compared to alpha contribution from excited and ground states. Hence, continuum states show similar behavior of wetting layer in the calculation model without continuum states.

• However, broadband QD lasers show that alpha contribution from continuum states is the largest and affect the alpha value of the system.

Z. Mi et al., IEEE J. Quantum. Electron., 43, 5 (2007).

EHIGH IVERSITY SIMULATION RESULTS (con't) Technologies



• The value of differential gain of typical QD lasers (Type I) shows similar magnitude of order as reported so far.

• As inhomogeneity of the system increase, differential gain at threshold decrease as compared to typical QD lasers.

• However, the decrease of differential gain is less than 1×10^{-16} cm².

• Hence, the degree of modulation bandwidth of broadband lasers will not be degraded too much from typical QD lasers.





Conclusions

• Derivative characteristics of broadband QD lasers are compared with typical QD lasers.

• A slight increase of linewidth enhancement factor at ground state of 0.8 is observed in broadband QD lasers.

• A slight decrease of differential gain of les than 1x10⁻¹⁶ cm² at threshold is observed in broadband QD lasers.

• The comparable values of GS's LEF in broadband QD lasers shows its competency in providing low frequency chirping as well as a platform of monolithic integration for the GS operation.





Thank You





Simulation parameters

E _{GS} = 1050meV	T _{0,wu} = 1ps
E _{ES} = 1090meV	$T_{0,ue} = T_{0,eg} = T_{0,ug} = 3.4ps$
D _G = 1	n _r = 3.5
D _E = 3	T _r = 1ns
D _U = 10	$N_{\rm D} = 1.67^* 10^{23}$
L _{ca} = 800 μ m	$V_A = 9.6^* 10^{-16}$
$R_1 = R_2 = 0.3$	$B = 10^{-4}$
$\alpha_{i} = 4.5 \text{ cm}^{-1}$	Γ _{QD} = 0.03
$T_{wr} = 0.4$ ns	Γ _{WL} = 0.01
T _e = 0.38ns	$\Gamma_{inhomo} = 23 meV, 76 meV$
$T_{0,uw} = 10ps$	△ E = 0.22meV