Dynamic Internal Optical Field Patterns for Self-Pulsating Two-Section DFB Lasers

Jer-Shien Chen¹, Hong-Chang Kung²,

San-Liang Lee³, and Hen-Wai Tsao¹

 ¹Graduate Institute of Communication Engineering and Department of Electrical Engineering, National Taiwan University
²Electrical Engineering Department of Tung Nan Institute of Technology
³Graduate Institute of Electronic Engineering and Department of Electronic Engineering, National Taiwan University of Science and Technology

1



Outline

- Introduction
- Existed Simulation Tools
- Optical Field Analysis and Numerical Calculation & Three Cases
- Conclusion
- Reference
 - Appendix: inte Model Math Equations

output waveform vs. internal optical field pattern

Introduction

- Investigate the mode beating self-pulsation in a two-section distributed feedback (TS-DFB) laser.
- Study the relationship between output waveform and optical field patterns for with/without AR-coating cases.
- It is more difficult to measure the internal variation inside the device, especial for the internal optical field analysis in space distribution and in time domain.
- What is the relationship between output waveform and internal optical field pattern?

Existed Simulation Tools

OptiSystem

- Optical communication system
- Providing laser module
- No self-pulsating simulation
- OptiSim
 - VPItransmissionMaker WDM VPIcomponentMaker Active Photonics
 - Optical communication system
 - Providing laser module with detailed parameters
 - Transmission Line Laser Model
 - Providing self-pulsating simulation output waveform

- 1. Finite-Difference method works well for laser, with spectral mixing, or with noise.
- The typical field patterns for the DFB with λ/4 phase shift have a peak on the location of phase shift.
- The finite difference method could be also applied to analyze the multi-section DFB. In addition, setting different section length will create mode depression.

- 4. In a clock source, the excitation is due to random 'spontaneous noise'. Hence, the outputs are not precisely the same for run-to-run, or iterations.
- 5. The result traces of iteration-to-iteration could be calculated, and the accumulative field pattern of all iterations could be visualized.
- 6. It is interesting to view the difference between optical field under non-selfpulsation and optical field under selfpulsation.

7. For the former, the **output is stable**. Hence, it is intuitional to image that the optical field should be more stable. That is, for different calculation iterations, the optical field patterns in space should be "almost over-lapped (Case1)".

7

8. For the latter, we consider two branches: anti-reflective (AR) coating and natural facets.





Calculated normalized space optical field (normalized relative power) for the proposed TS-DFB without self-pulsation, including forward field (from left to right, the leftest point has smallest field), reverse field (from left to right, the leftest point has largest field), and sum field.

9. The internal optical field patterns should be dynamic for each periodic waveform; that is, it is difficult for the optical field patterns in space to be almost over-lapped. Instead, it is changed (Case2).

 A TS-DFB may be processed with AR coating on both facets. Otherwise, the reflective effect from natural facets should be considered (Case3).



 Calculated normalized optical power spectrum for TS-DFB (with ARcoating) mode-beating mechanism under self-pulsation for long time simulation. The central wavelength is 1550nm.



Calculated normalized space optical sum field (normalized relative power) for TS-DFB including AR-coating, with mode-beating mechanism, and under self-pulsation for long time simulation. Two optical fields for "mountain-peak-swinging" could be observed.



for TS-DFB mode-beating mechanism, under self-pulsation for long time simulation. The pulsation period is 25ps (40GHz).

Calculated normalized optical power spectrum for TS-DFB (with natural facets) under self-pulsation for long time simulation.

Calculated normalized space optical sum field (normalized relative power) for the TS-DFB including natural facets, under self-pulsation for long time simulation. Mixing "almost over-lapped" and "mountain-peak-swinging" could be observed.

Conclusion

- We first apply the Finite-Difference method to DFB internal optical field patterns to study the stable output case.
- Then we analyze the internal optical field patterns of TSDFB under self-pulsation; and our results show that the field patterns also vibrate between two peaks from run-to-run.
- Finally, the difference effect between ARcoating and the natural facets also has been studied.
- In short, the relationship between output waveform and internal optical field pattern has been studied.

Reference

- Sartorius.B, et.al,."Dispersive self-Qswitching in self-pulsating DFB lasers," IEEE Journal of Quantum Electronics, Vol.33, No.2, pp.211-218, Feb 1998.
- Marcenac.D.D, et.al,."Distinction between multimoded and singlemoded self-pulsationsin DFB lasers," IEEE Electronics Letters, Vol.30, No.14, pp.1137-1138, July 1994.
- John Carroll, et. al.. 'Distributed Feedback semiconductor lasers', chap.7 (1998).

Appendix: Model Math Equations

$$\frac{1}{v_g}\frac{\partial F}{\partial t} + \frac{\partial F}{\partial z} = j\kappa R + (g - j\delta)F + i_{spf}$$
$$\frac{1}{v_g}\frac{\partial R}{\partial t} - \frac{\partial R}{\partial z} = j\kappa F + (g - j\delta)R + i_{spr}$$

$$\begin{bmatrix} F\{(T+1), (Z+1)\}\\ R\{(T+1), Z\} \end{bmatrix} = \exp\{(g-j\delta)s\} \begin{bmatrix} \cos\theta & j\sin\theta\\ j\sin\theta & \cos\theta \end{bmatrix} \begin{bmatrix} F\{(T,Z)\}\\ R\{T, (Z+1)\} \end{bmatrix}$$
$$\approx \left(\frac{2+gs-j\delta s}{2-gs+j\delta s}\right) \begin{bmatrix} \sqrt{1-\kappa^2 s^2} & j\kappa s\\ j\kappa s & \sqrt{1-\kappa^2 s^2} \end{bmatrix} \begin{bmatrix} F\{(T,Z)\}\\ R\{T, (Z+1)\} \end{bmatrix}$$
$$\sin\theta = \frac{\kappa s}{1+\frac{1}{4}\kappa^2 s^2} \approx \kappa s \quad \cos\theta \approx \sqrt{1-\kappa^2 s^2}$$