



Columnar quantum dots (QD) in polarization insensitive SOA and non-radiative Auger processes in QD: a theoretical study

J. Even, L. Pedesseau, F. Doré, *S. Boyer-Richard*, UMR FOTON 6082 CNRS, INSA de Rennes, France Jacky.even@insa-rennes.fr



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### Quantum dot with axial symmetry



Geometry of quantum dot :  $C_{\infty v}$  symmetry



A new description in cylindrical coordinates :  $\mathbf{r}, \boldsymbol{\varphi}, \mathbf{z}$ 

Some references for k.p axial approximation :

<sup>1</sup> P. Enders and M. Woerner, SST (1996) : k-dependent 4x4 block-diag. of 8x8 Ham. (Bulk)

<sup>2</sup> C. Y. P. Chao and S. L. Chuang, PRB (1992) : k-ind. block-diag. of 6x6 Ham. (QW)

<sup>3</sup> Y. M. Mu and S. S. Pei, JAP (2004) : k-ind. block-diag. of 8x8 Ham. (QW)

<sup>4</sup>K. J. Vahala and P. C. Sercel, PRL, PRB (1990) : k-ind. block-diag. of 8x8 Ham. (Qwire and Spherical QD)

<sup>5</sup> M. Tadic, F.M. Peeters and K. J. Janssens, JAP, PRB (2002, 2004) k-ind. block-diag. of 6x6 Ham. (QD)



# Axial approximation for mechanical properties



Acoustic phonon **band "warping**" in cubic materials







#### Axial approximation for the 8-band k.p strained Hamiltonian



**Unstrained part :**  $R = -\sqrt{3} \frac{\hbar^2}{2m_0} \left[ \gamma_2 \left( k_x^2 - k_y^2 \right) - 2i \gamma_3 k_x k_y \right] \approx -\sqrt{3} \frac{\hbar^2}{2m_0} \overline{\gamma} k_-^2$  $\boldsymbol{R}_{\boldsymbol{\varepsilon}} = \frac{b\sqrt{3}}{2} \left(\boldsymbol{\varepsilon}_{rr} - \boldsymbol{\varepsilon}_{\varphi\varphi}\right) \cos(2\boldsymbol{\varphi}) - i\frac{d}{2} \left(\boldsymbol{\varepsilon}_{rr} - \boldsymbol{\varepsilon}_{\varphi\varphi}\right) \sin(2\boldsymbol{\varphi}) \approx \frac{\overline{b}\sqrt{3}}{2} \left(\boldsymbol{\varepsilon}_{rr} - \boldsymbol{\varepsilon}_{\varphi\varphi}\right) e^{-i2\varphi}$ Strained part : (new proposition)  $\frac{\overline{b}\sqrt{3}}{2} = \frac{1}{2} \left( \frac{b\sqrt{3}}{2} + \frac{d}{2} \right) \begin{cases} \text{InAs} & \frac{b\sqrt{3}}{2} = -1.58eV & \frac{d}{2} = -1.80eV \\ \text{GaAs} & \frac{b\sqrt{3}}{2} = -1.56eV & \frac{d}{2} = -2.25eV \\ \text{InP} & \frac{b\sqrt{3}}{2} = -1.73eV & \frac{d}{2} = -2.50eV \end{cases}$ Even et al, PRB (2008)  $Q_{\varepsilon} = b \left( \varepsilon_{zz} - \frac{\varepsilon_{rr} + \varepsilon_{\varphi\varphi}}{2} \right) \qquad R_{\varepsilon} = \frac{\overline{b} \sqrt{3}}{2} \left( \varepsilon_{rr} - \varepsilon_{\varphi\varphi} \right) e^{-i2\varphi}$  $S_{\varepsilon} = -d\varepsilon_{rz} e^{-i\varphi}$  $A_{\varepsilon} = a_{c} \left( \varepsilon_{rr} + \varepsilon_{\varphi\varphi} + \varepsilon_{zz} \right)$  $\boldsymbol{P}_{\varepsilon} = \boldsymbol{a}_{v} \left( \boldsymbol{\varepsilon}_{rr} + \boldsymbol{\varepsilon}_{\boldsymbol{\omega}\boldsymbol{\omega}} + \boldsymbol{\varepsilon}_{\boldsymbol{\tau}\boldsymbol{\tau}} \right)$ Hydrostatic strain **Biaxial strain** Shear strain



#### Axial approximation for the 8-band k.p strained Hamiltonian



#### Block diagonalization of the Hamiltonian for each Fz value :

A good quantum number : total angular momentum  $F_z = J_z + L_z$ 

Development of the wavefunction :  $|J, J_z\rangle |L_z = F_z - J_z\rangle$ 

Basis of 8 Bloch functions  $u_i$  (i=1...8) 8 enveloppe functions of (r,z)

| J <sub>Z</sub> | Fz              | -5/2    | -3/2 | -1/2 | 1/2 | 3/2 | 5/2 |
|----------------|-----------------|---------|------|------|-----|-----|-----|
| 1/2            | L <sub>z1</sub> | -2<br>R | -1   | 0    | +1  | +2  | +3  |
| -1/2           | L <sub>z2</sub> | -3      | -2   | -1   | 0   | +1  | +2  |
| -1/2           | L <sub>z3</sub> | -3      | -2   | -1   | 0   | +1  | +2  |
| -3/2           | L <sub>z4</sub> | -4      | -3   | -2   | -1  | 0   | +1  |
| 3/2            | L <sub>z5</sub> | -1      | 0    | +1   | +2  | +3  | +4  |
| 1/2            | L <sub>z6</sub> | -2      | -1   | 0    | +1  | +2  | +3  |
| 1/2            | $L_{z7}$        | -2      | -1   | 0    | +1  | +2  | +3  |
| -1/2           | $L_{z8}$        | -3      | -2   | -1   | 0   | +1  | +2  |



Axial approximation for the 8-band k.p strained Hamiltonian



#### **Excited "p"and "d" states splitting predicted** from symmetry analysis













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Auger processes in narrow gap QD





Example : InAs/Q1.18/InP truncated cone QD  $E_{CH}=0.77 eV (gap) E_{CC}=23 meV E_{HH}=16 meV$  $F_{z}=1/2$   $N_{Bulk}=10^{18} \text{ cm}^{-3}$ Non-radiative processes :  $\tau_{\rm CHCC} = 0.2 \mu s \tau_{\rm CHLH} = 1.1 m s (\tau_{\rm CHSH} = 14 n s)$ CB relaxation in QD :  $\tau_{\rm CCCC} = 0.74 \, {\rm ps} \, \tau_{\rm CCLH} = 24 \, {\rm ps} \, \tau_{\rm CCSH} = \infty$ Hole relaxation in QD :  $\tau_{\rm HHCC}$ =1.0ps)  $\tau_{\rm HHLH}$ =22ps  $\tau_{\rm HHSH}$ =  $\infty$ 





## Conclusion : further studies...

- New axially symmetric strained nanostructures
- Auger effects (gap influence, comparison WL/bulk, barrier materials, hydrostatic pressure...)
- Beyond the 8-band k.p approximation

Soline Boyer-Richard, NUSOD'08, Nottingham, 4th September 2008