

#### THALES

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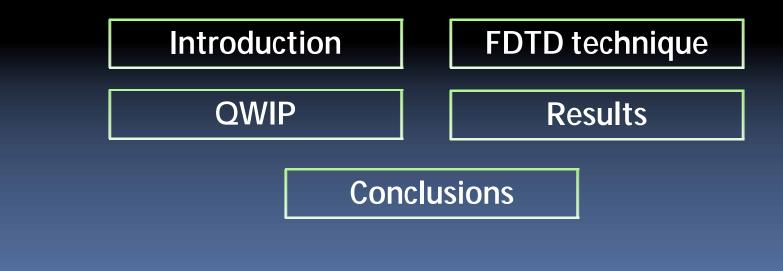
#### ACCURATE SIMULATION OF QUANTUM WELL INFRARED PHOTODETECTORS BY FDTD



#### Aim of the Presentation

To show our approach based on the Finite Difference Time Domain (FDTD) technique for the design and the optimization of Quantum Well Infrared Photodetector (QWIP) devices

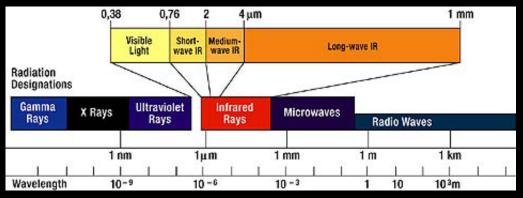
Presentation Outline



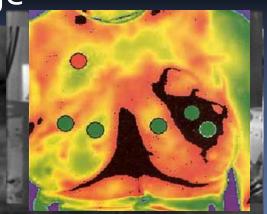


### Introduction

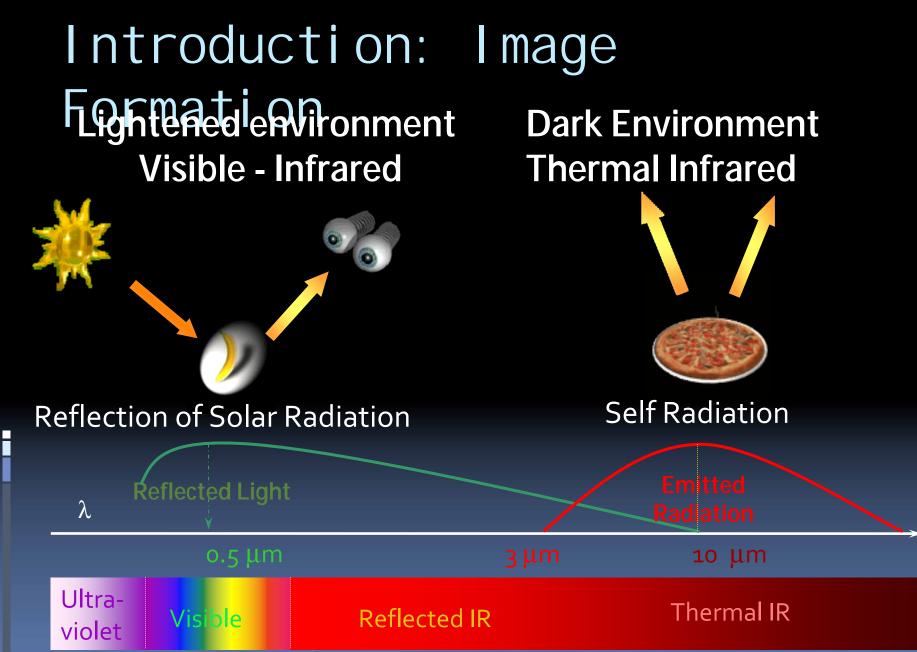
The Infrared Domain: 0.76 μm ÷ 1000 μm



- Efficient Technology useful to detect the e.m. radiation in the IR Spectrum Range
- Possible applications:
  - Civil (Security, Surveillance)
  - Medical (Brest Cancer Detection)
  - Military









### Introduction

- The Image Formation
  - Lightened environment:
    - Detection of the photons emitted by a light source and reflected by the objects
  - Dark environment:
    - Objects at non-zero temperature emit photons

Planck's law 
$$E(\lambda,T) = \frac{2hc^2}{\lambda^5} \frac{1}{e^{\frac{hc}{\lambda Tk}} - 1} \begin{bmatrix} W \\ m^2 \mu m \end{bmatrix} \begin{pmatrix} c = 310^{-8} m/s \\ h = 6.62510^{-27} erg/s \\ k = 1.3810^{-27} erg/^{\circ}K$$

 Using different detectors it is possible to build up different images of the same scenario



#### Introduction

 Images of the same human head obtained by different detection techniques



Color Image Visible Band 0.38 µm ÷ 0.76 µm



B&W Image Visible Band 0.38 μm ÷ 0.76 μm



Solar Reflection IR Band 1.0 µm ÷ 1.7 µm



Solar Reflection + Thermal Emission IR Band 3.4  $\mu$ m ÷ 5.0  $\mu$ m



Thermal Emission IR Band 8.5  $\mu$ m ÷ 9.5  $\mu$ m



#### Infrared Detectors

#### PhotoraicDetectors

- Absorption of IR radiation (\* In ErgerseEdectronthe detector
   Absorption of IR radiation (\* In ErgerseEdectronthe detector
   Absorption of IR radiation (\* In ErgerseEdectronthe detector)
- $\Rightarrow$  Only  $\lambda$ :  $E_g > E_{go}$  can be absorbed
- 8 Problems in obtaining detectors for long wave ( $\lambda \approx 10 \ \mu m$ ) IR (small E<sub>g</sub> materials:  $E_g \approx 0.1 \ eV$ )
- Conventional Detection' (with weak band-gap materials as Hg<sub>x</sub>-Cd<sub>1-x</sub>Te) is not efficient (exotic materials, not developed!)
- `Effective' band-gap materials (GaAs/AlGaAs heterostructures) which use InterSubBand transitions created by Quantum Wells in large-band-gap semiconductors



#### Quantum Well Infrared Detectors Photonic detectors with weak band-gap (Hg<sub>x</sub>-Cd<sub>1-x</sub>Te)

- (to have detection @ long wave IR)
- $\boldsymbol{X}$  concentration adjusted for  $\lambda$  tuning
- Electrons excited from VB to CB with inter-sub-band transitions
- Not very well developed (exotic materials!)

 Quantum Well detectors based on GaAs/AlGaAs heterostructures

- Developed since 1990
- Multi Quantum Well Photonic detectors (QWIP)

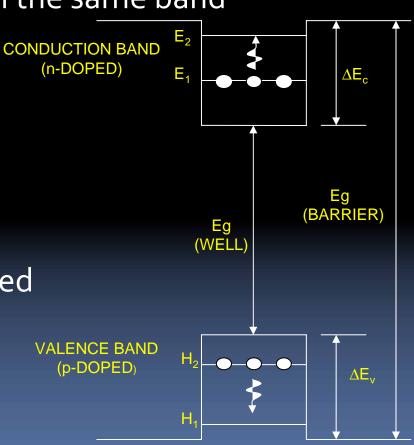
#### Multi Quantum Well IR Detectors

- The Detection Process: Inter-Sub-Band transitions
  - Involves transitions within the same band

Augumentum Well needed

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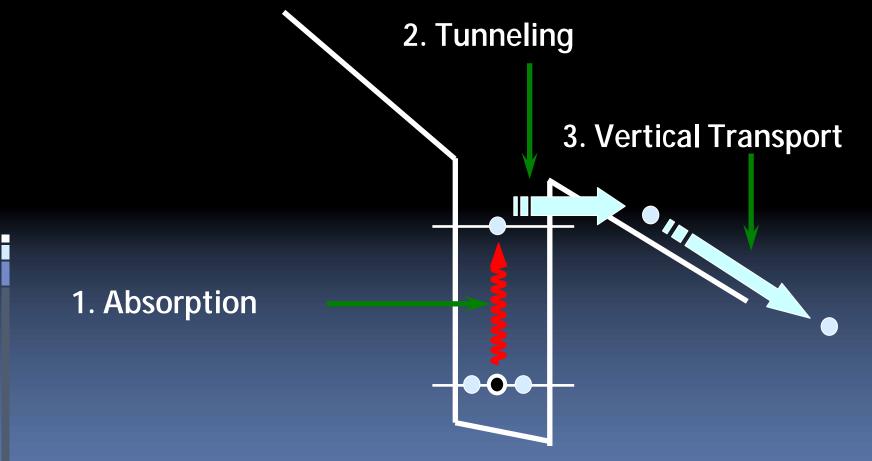
 Electron (Hole) from the doped QW ground state in CB (VB) to un unoccupied state in the same band





#### Multi Quantum Well IR Detectors

 QW structure designed to have carrier escaping from the well and collected as a (photo) current



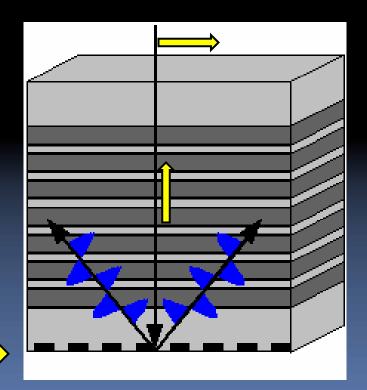
#### Multi Quantum Well IR Detectors

#### InterSubBand Transitions

 Energy levels inside CB or VB arise from the spatial localization introduced in the QW of a low-band-gap material (GaAs) surrounded by a higher-band-gap semiconductor (Al<sub>x</sub>Ga<sub>1-x</sub>As)

#### Optical Absorption:

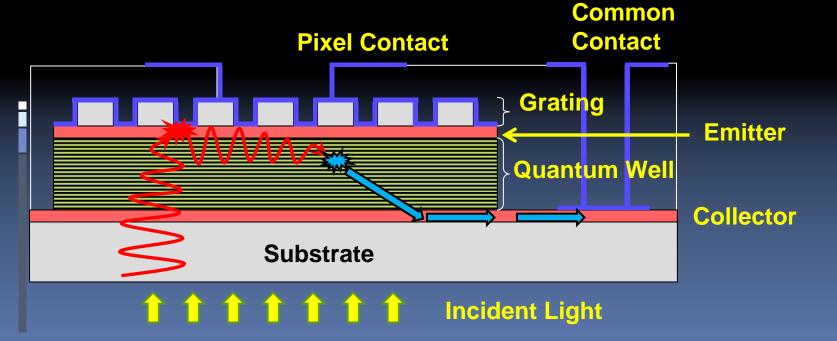
- only the optical field along the superlattice direction (made by the well-barrier structure) is absorbed
- Light (TEM polarized) orthogonally polarized respect to the direction of interest
- ➡ Polarization rotation (TEM ➡ TM) is needed!
- ➡ Diffraction gratings are used



#### Quantum Well IR Photodetector

- The Structure of a QW IR Photodetector
  - Substrate: GaAs

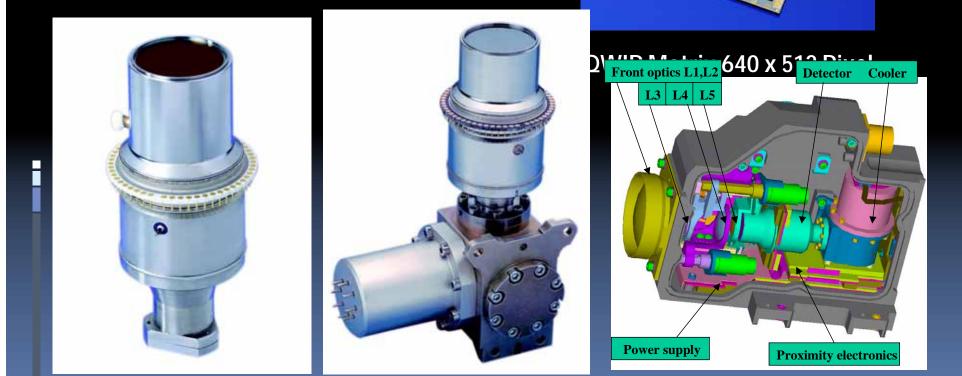
- Collector and Emitter: doped GaAs:Si
- Active zone: 40 QWs (doped GaAs, barrier: AlGaAs)
- Grating: GaAs + metallic coat (Au, Ni)





#### Thermal Imager (IR Camera)

- Matciceseteeteenapet in the detector system agginstalled inside the IR Camera
  - Each Pixel is a QWIP





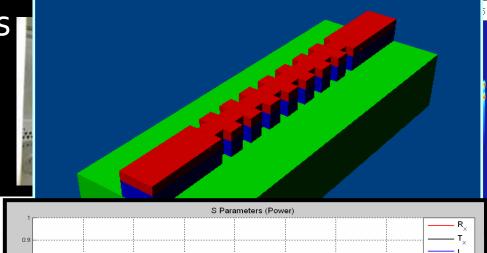
### The FDTD Approach

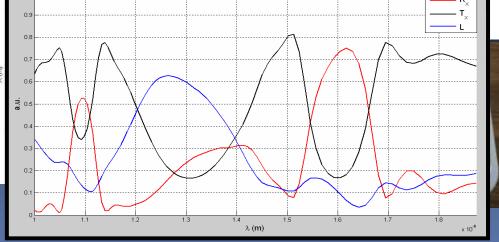
 The Finite Difference in the Time Domain (FDTD) approach is used to design and optimize the performance of QWIPs

#### Why FDTD?

- Available in our group
- FDTD Properties

Microwave Heating Optics Applications







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### The FDTD Approach

#### The FDTD Properties

- Generality and Versatility
- Dissipative, Dispersive and Non Linear materials can be 'easily' included
- Temporal evolution of the e.m. fields
- Frequency Domain results available by Fourier Transform

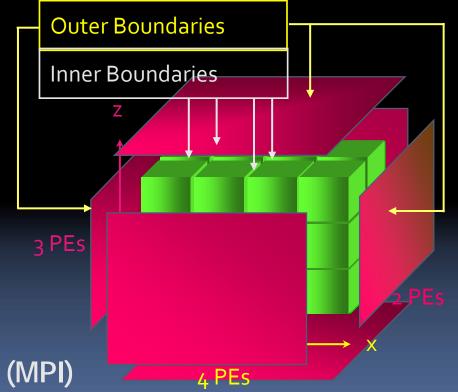
#### <sup>®</sup> Time and Memory consuming

- ${\ensuremath{^{&}}} {\ensuremath{^{&}}} Problems with devices of several <math display="inline">\lambda$  on each side are impractical on a simple PC
  - Lack of Computer Memory (RAM)
  - Long CPU time
- ⇒ Use of Parallel Computing



### Parallel FDTD Technique

- FDTD is well suited for Parallel Computation as the solving algorithm mainly involves 'local data'
- Domain Decomposition
   Each Block belongs to a single PE
- Boundary Conditions
  - Outer Boundaries
     ABC PML 3PE
     Inner Boundaries
     Data Communication
     Message Passing Interface (MPI)

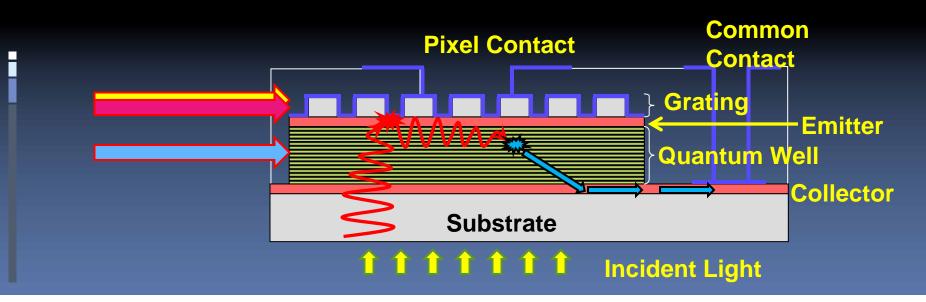


Tot = 24 PEs



### FDTD Simulation of QWIPs

- Simulation Strategy:
  - Optimization of the Metallic Grating
    - PEC metallic surface
    - Real metal (Drude Model)
  - Lorentz Model for the InterSubBand Absorption
  - TFSF Approach for Field Excitation





### Grating Optimization

- Coupling grating fundamental component of a QWIP (only TM waves are absorbed, but the incident light is mainly TEM Polarized)
  - TEM to TM Polarization Rotation
  - Increasing of the e.m. field in the active region (Surface Plasmons + Surface Cavity Effect)

# ⇒Grating optimization is essential for good QWIP performance

L. Stabellini, M. Carras, A. De Rossi and G. Bellanca, "Design and Optimization of High-Q Surface ModeCavities on Patterned Metallic Surfaces", IEEE JQE, 2008, in press



### Real Conductor

 Drude model used to describe the interaction between the light and the 'real' conductor

$$\varepsilon_{0} \frac{\partial \vec{E}}{\partial t} = \nabla \times \vec{H} - \vec{J}_{p}$$

$$\frac{\partial \vec{J}_{p}}{\partial t} = -\upsilon \vec{J}_{p} + \varepsilon_{0} \omega_{p}^{2} \vec{E}$$

$$\hat{\varepsilon} = \varepsilon_{0} \left[ 1 + \frac{\omega_{p}^{2}}{\omega(j\upsilon - \omega)} \right] = \varepsilon_{0} [1 + \hat{\chi}(\omega)]$$

 $ω_p$ : Plasma frequency v: Collision frequency

 Implemented in FDTD using the Auxiliary Differential Equation (ADE) technique



### Lorentz Model for InterSubBand Absorption

 A Lorentz model can be used to describe InterSubBand absorption of a Quantum Well IR Photodetectors (A. Nedelcu, 'Detection Infrarouge, Imaginerie Infrarouge', Thales Internal Report)

$$\varepsilon_{0} \frac{\partial \vec{E}}{\partial t} = \nabla \times \vec{H} - \vec{J}_{p}$$

$$\frac{\partial \vec{J}_{p}}{\partial t} = -\upsilon_{1} \vec{J}_{p} + (\varepsilon_{s} - \varepsilon_{\infty}) \varepsilon_{0} \omega_{1}^{2} \vec{P} \quad \frac{\partial \vec{P}}{\partial t} = \vec{J}_{p}$$

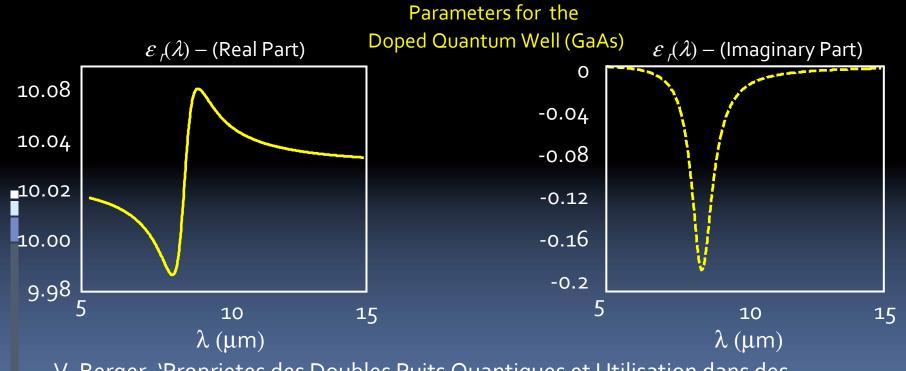
$$\hat{\varepsilon} = \varepsilon_{0} \varepsilon_{\infty} + \frac{(\varepsilon_{s} - \varepsilon_{\infty}) \varepsilon_{0} \omega_{1}^{2}}{\omega_{1}^{2} + j \omega \upsilon_{1} - \omega^{2}} = \varepsilon_{0} [1 + \hat{\chi}(\omega)]$$

- $\omega_1$ : Resonant frequency
- $\upsilon_1$ : Damping frequency
- $\varepsilon_s$ : Static relative permittivity
- $\varepsilon_{\infty}$ : Infinite relative permittivity
- Implemented in FDTD using the ADE Technique
   Same model for both Drude and Lorentz media
- Multi-Pole Lorentz model integrates the two different material representations in a single procedure (Drude material = zero order pole)



#### Parameters for the Lorentz

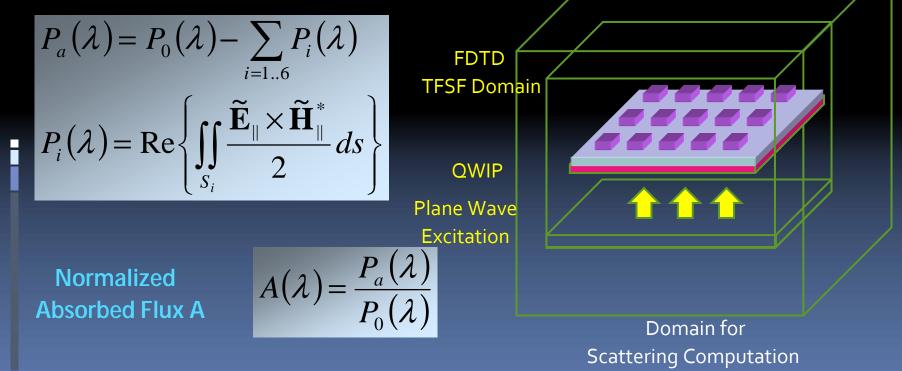
**Note** arameters  $\omega_1$ ,  $\upsilon_1$ ,  $\varepsilon_\infty$  and  $\Delta \varepsilon = \varepsilon_s - \varepsilon_\infty$  can be obtained starting from the Density Matrix formalism and considering the doping parameters and the refractive index of the semiconductors used in the active region



V. Berger, 'Proprietes des Doubles Puits Quantiques et Utilisation dans des Dispositif Optoelectroniques', PhD Thesis, Paris VI, 1992)

### Spectral Response of a QWIP

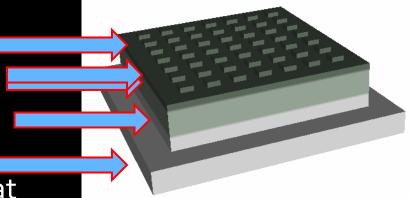
- Excitation of a 'Plane Wave' propagating in y direction
- Computation of the field 'Scattered' by the QWIP
- Computation of the field 'absorbed' by the semiconductor

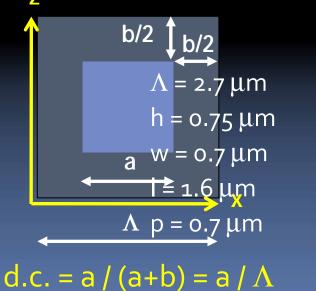


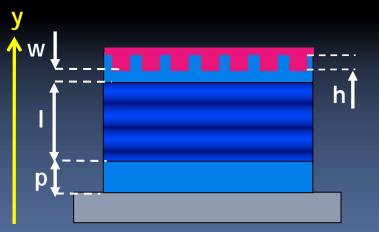


### QWIP Parameters

- Complete simulated structure for a realistic QWIP
  - Substrate
  - GaAs:Si Collector
  - Active Zone: 40 QW (1.6μm)
  - GaAs:Si Emitter
  - Grating (GaAs) + metallic coat







Grating Emitter QW Region Collector Substrate

### FDTD Simulation Parameters

#### Discretization

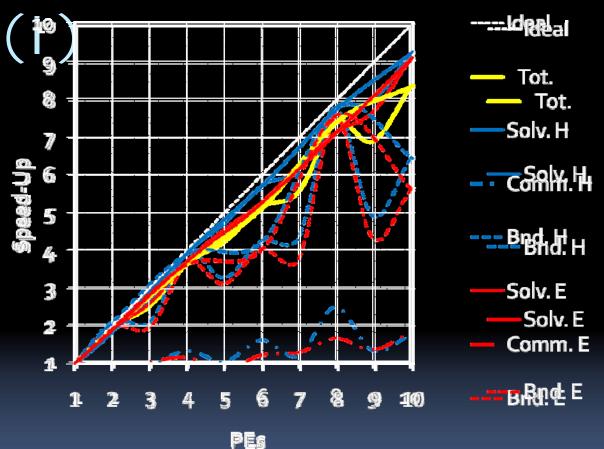
- $\Delta_x = \Delta_y = \Delta_z = 75 \text{ nm} (36 \text{ points}/\Lambda)$
- 7Λ x 7Λ QWIP: 277 x 63 x 277 cells
- PML Layer: 14 cells, ρ = 1.0 e<sup>-6</sup> (8,465,275 overall mesh points)
- Temporal Time Step: 7.22 e<sup>-16</sup> s

#### Input / Output Parameters

- TFSF Excitation:  $f_0 = 35$  THz ( $\lambda = 8.5 \mu m$ ); BW = 12THz
- DFT Computation: frequency range [25÷45] THz (50 samples)
- Number of Time steps: 30000
- Total Computation Time: 8600 s (~ 2.4 h on 6 PEs PIV 3GHz)



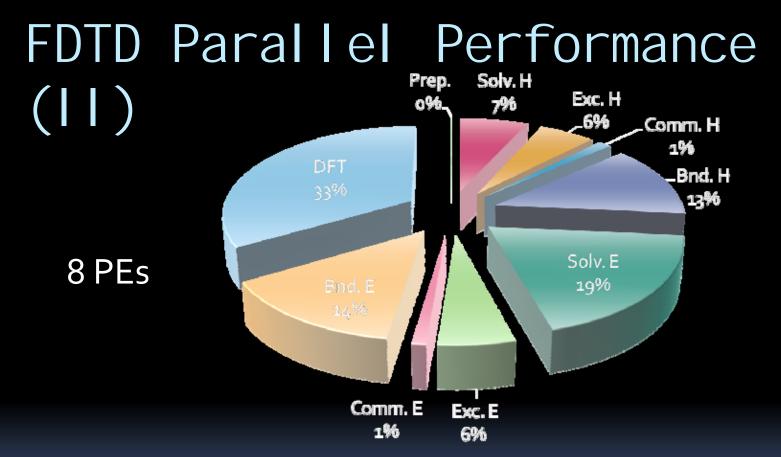
#### FDTD Parallel Performance



- G889 scaling of pure computations' (Yee's Solver: Solv. H, Solv. E) Boundary Gonditions dep't scale as good as the Yee's Solver
  - Η
  - No good performance with an odd number of PEs; 8.3 with 10 PEs H



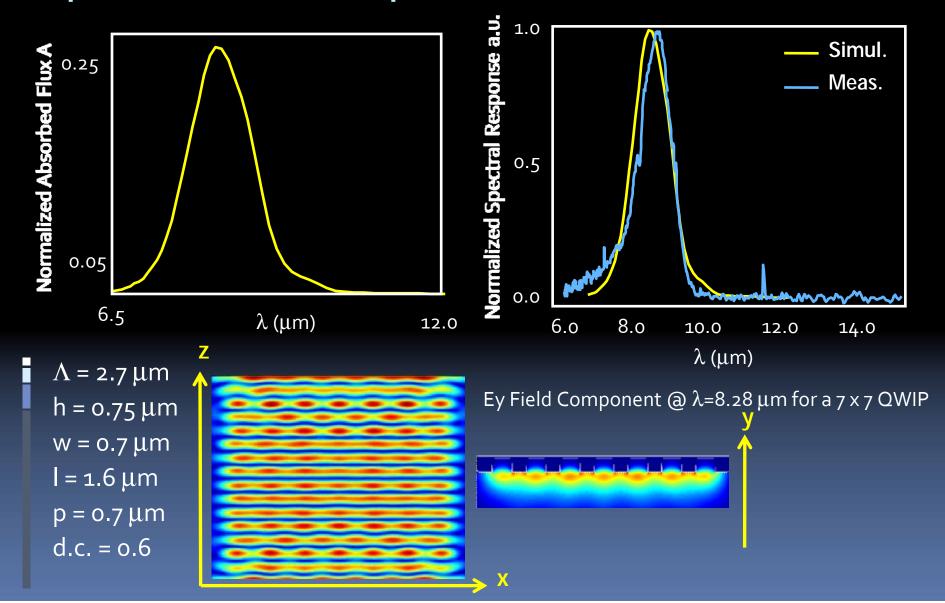
H



- E Solver more time consuming than H Solver because of J and P computations (Dispersive Materials in the QW region)
- Good work balance between E and H in Communication, Excitation and Boundary Computation
- DFT computationally intensive



#### Spectral Response of a QWIP





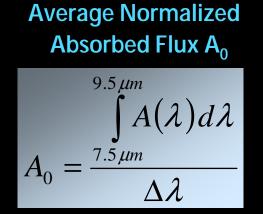
#### **QWIP** Optimization 7 × 7 structure (20.5 × 20.5 µm) 0.18 $\Lambda$ = 2.7 $\mu$ m **Average Normalized** Absorbed Flux A<sub>0</sub> h = 0<del>.7</del>5.6mm $w = 0.7 \,\mu m$ 9.5 µm $(\lambda)d\lambda$ l = 1.6 μm $p = 0.7 \,\mu m$ $7.5\,\mu m$ 0.20 d.c. = 0.6 0.02 PEC 0.8 1.6 Corrugation Depth (µm) Optimum value for the **PEC** good approximation verage Normalize Corrugation Depth: $h=0.75 \,\mu m$ for the metallic coating (Au) 0.02 1.6 0 Corrugation Depth (µm)



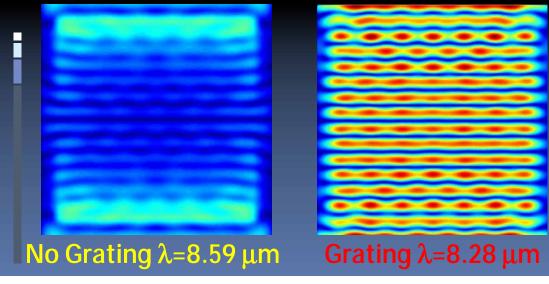
# QWIP Optimization (II)

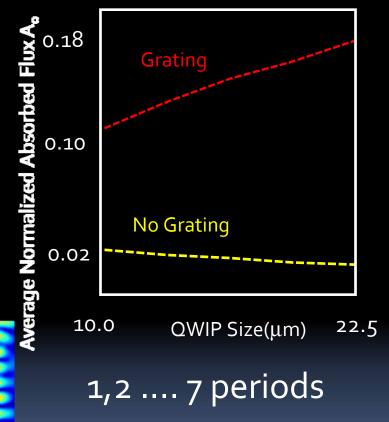
#### 7 × 7 structure (20.5 × 20.5 μm)

Λ = 2.7 μm h = 0.75 μm w = 0.7 μm I = 1.6 μm p = 0.7 μmd.c. = 0.6



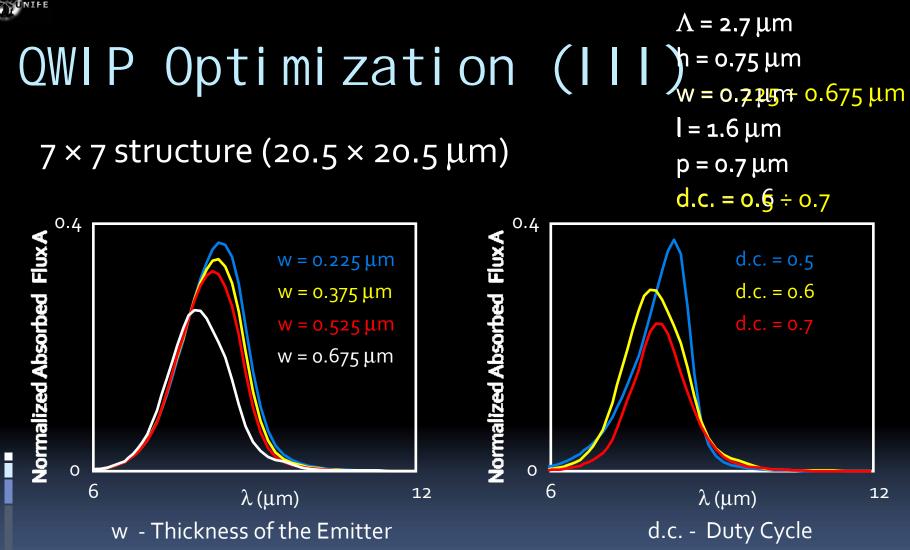
Ey Field Component (a)  $\lambda$  of the maximum absorption





Metallic grating is fundamental for optimum performance of the QWIP device





Reducing the thickness of the emitters w, the distance between the grating and the active zone decreases, thus increasing the E<sub>y</sub> field in the QW zone
 d.c. = 0.5 allows the best performance of the QWIP



## 1D Coupling Grating QWIP

- 1D Coupling Grating used for 'polarization sensitive' devigesperimental Results
  - Only one linear component of the generally elliptically polarized incident light should be detected

 Detection of images with low thermal contrast or cluttered scenes

 <sup>E</sup>
 <sup>L</sup>
 <sup></sup>

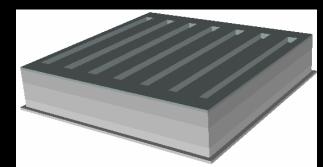
 Combining signals from pixels of 1D gratings oriented differently, the full characterization of a linear polarization degree in a scene is allowed

6.0 6.5 7.0 7.5 8.0 8.5 9.0 9.5 10.0

λ (μm)

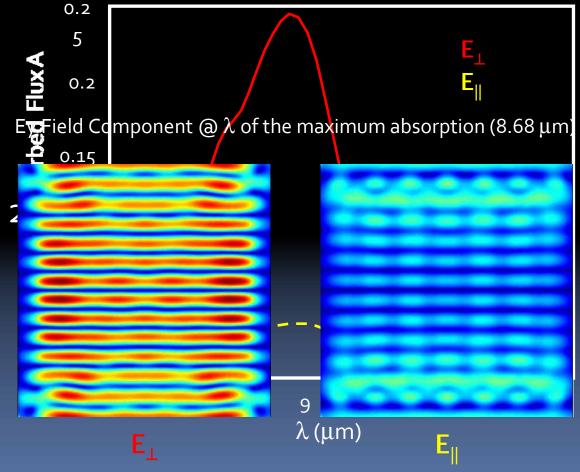


# 1D Coupling Grating -Results



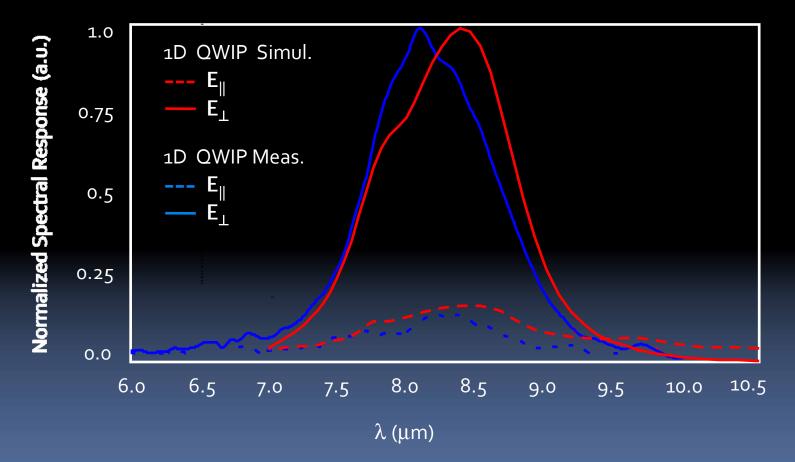
#### 7 × 1 structure (20.5 × 2

 $Λ = 2.7 \mu m$   $h = 0.75 \mu m$   $w = 0.7 \mu m$   $l = 1.6 \mu m$   $p = 0.7 \mu m$ d.c. = 0.5





# 1D Coupling Grating -Resul **t**imulation vs Measurement





#### Conclusions

- FDTD technique used as a design and optimization tool for QWIP devices
  - Design and Optimization of the Grating Surface
  - Investigations on the influence of the different parameters on the absorption of a QWIP
  - 2D and 1D coupling grating investigated
  - Good agreement between simulations and measurements