### Many-body optical gain in ZnO- and GaNbased quantum well lasers

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#### Overview

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(0001)-oriented QW structures
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Summary



### Introduction

#### Wurtzite GaN-based QW Laser

Potential and existing optoelectronic device applications : laser diode, traffic lights, displays, and so forth

- Properties of (0001)-oriented WZ GaN-based QW lasers: Several disadvantages, compared to conventional ZB GaAs- or InP-based QW lasers
- 1) GaN : significantly larger effective mass
- 2) Biaxial strain : does not effectively reduce effective mass
- 3) Large internal field : PZ and SP polarizations

→ Higher threshold current density for lasing
 [ Control the internal field and the effective mass ]



### Main Goals of the Approach

- The crystal orientation effects : a new parameter in the band structure engineering
  - ← by controlling the internal field and the effective masses
- Recently, ZnO and related oxides :

   new wide band-gap semiconductors
- 1) Growth temperature of ZnO : around 500 °C (GaN: 1000 °C)
- 2) The internal field is expected to be much smaller than that in GaN-related systems.

Electronic and optical properties of ZnO/MgZnO QW structures and Crystal orientation effects

Compare with those of GaN-based QW structures



### **Theoretical background**





### Valence band structure

#### (0001)-oriented Hamiltonian

$$H(\mathbf{k},\epsilon) = \begin{pmatrix} F & -K^* & -H^* & 0 & 0 & 0 \\ -K & G & H & 0 & 0 & \Delta \\ -H & H^* & \lambda & 0 & \Delta & 0 \\ 0 & 0 & 0 & F & -K & H \\ 0 & 0 & \Delta & -K^* & G & -H^* \\ 0 & \Delta & 0 & H^* & -H & \lambda \end{pmatrix} \begin{vmatrix} U_1 \\ U_2 \\ U_3 \\ U_4 \\ U_5 \\ U_6 \end{pmatrix} = \begin{vmatrix} F & \Delta_1 + \Delta_2 + \lambda + \theta, \\ G &= \Delta_1 - \Delta_2 + \lambda + \theta, \\ A &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \lambda_{\epsilon}, \\ B &= \frac{\hbar^2}{2m_o} [A_3k_z^2 + A_4(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ H &= \frac{\hbar^2}{2m_o} [A_3k_z^2 + A_4(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ K &= \frac{\hbar^2}{2m_o} A_5(k_x + ik_y)^2 + D_5\epsilon_+, \\ H &= \frac{\hbar^2}{2m_o} A_6(k_x + ik_y)(k_z) + D_6\epsilon_{z+}, \\ \lambda_{\epsilon} &= D_1(\epsilon_{zz}) + D_2(\epsilon_{xx} + \epsilon_{yy}), \\ \theta_{\epsilon} &= D_3(\epsilon_{zz}) + D_4(\epsilon_{xx} + \epsilon_{yy}), \\ \theta_{\epsilon} &= B_3(\epsilon_{zz}) + D_4(\epsilon_{xx} + \epsilon_{yy}), \\ \theta_{\epsilon} &= A_1 - \Delta_2 + \lambda + \theta, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \lambda_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_3k_z^2 + A_4(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ K &= \frac{\hbar^2}{2m_o} A_6(k_x + ik_y)(k_z) + D_6\epsilon_{z+}, \\ \lambda_{\epsilon} &= D_1(\epsilon_{zz}) + D_2(\epsilon_{xx} + \epsilon_{yy}), \\ \theta_{\epsilon} &= D_3(\epsilon_{zz}) + D_4(\epsilon_{xx} + \epsilon_{yy}), \\ \theta_{\epsilon} &= A_1 - \Delta_2 + \lambda + \theta, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \lambda_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z^2 + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_z + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_y + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_y + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_y + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_y + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}, \\ \lambda &= \frac{\hbar^2}{2m_o} [A_1k_y + A_2(k_x^2 + k_y^2)] + \theta_{\epsilon}$$



#### Rotation matrix

 $U = \begin{pmatrix} \cos\theta\cos\phi & \cos\theta\sin\phi & -\sin\theta \\ -\sin\phi & \cos\phi & 0 \\ \sin\theta\cos\phi & \sin\theta\sin\phi & \cos\theta \end{pmatrix}$ 



### $\varepsilon_{ij}(\theta) \leq \text{Minimization of strain energy}$



### Strain tensors in general coordinate



$$\epsilon_{xz} = -\frac{\left[(c_{11} + c_{12} + c_{13}\epsilon_{zz}^{(0)}/\epsilon_{xx}^{(0)})\sin^2\theta + (2c_{13} + c_{33}\epsilon_{zz}^{(0)}/\epsilon_{xx}^{(0)})\cos^2\theta\right]\epsilon_{xx}^{(0)}\cos\theta\sin\theta}{c_{11}\sin^4\theta + 2(c_{13} + 2c_{44})\sin^2\theta\cos^2\theta + c_{33}\cos^4\theta}$$

### Hamiltonian for general crystal orientation - Vectors and tensors



# PZ and SP polarizations for general crystal orientation

: Function of  $\theta$ 

 $P_{PZ} = P_x \sin \theta + P_z \cos \theta,$ 

 $P_{SP} = P_{SP}^{(0001)} \cos \theta$ 

 $P_x = 2d_{15}c_{44}\epsilon_{xz},$ 

 $P_z = [d_{31}(c_{11} + c_{12}) + d_{33}c_{13}](\epsilon_{xx} + \epsilon_{yy}) + [2d_{31}c_{13} + d_{33}c_{33}]\epsilon_{zz}$ 

$$P_x = 0$$
 for  $\theta = \pi/2$   
(  $\epsilon_{xz}=0$  )

(1010) Crystal orientation  

$$[\theta = \pi/2]$$
  
 $\Rightarrow P_{PZ} = 0$  and  $P_{SP} = 0$ 



## Non-Markovian gain model with many-body effects

$$g(\omega) = \sqrt{\frac{\mu_o}{\epsilon}} \left(\frac{e^2}{m_o^2 \omega}\right) \int_0^{2\pi} d\Phi \int_0^{\infty} dk_{||} \frac{2k_{||}}{(2\pi)^2 L_w} |M_{nm}(k_{||}, \Phi)|^2 [f_n^c(k_{||}, \Phi) - f_m^v(k_{||}, \Phi)] L(\omega, k_{||}, \Phi) = \frac{1}{2} \int_0^{2\pi} d\Phi \int_0^{\infty} dk_{||} \frac{2k_{||}}{(2\pi)^2 L_w} |M_{nm}(k_{||}, \Phi)|^2 [f_n^c(k_{||}, \Phi) - f_m^v(k_{||}, \Phi)] L(\omega, k_{||}, \Phi) = \frac{1}{2} \int_0^{2\pi} d\Phi \int_0^{2\pi} d\Phi \int_0^{2\pi} dk_{||} \frac{2k_{||}}{(2\pi)^2 L_w} |M_{nm}(k_{||}, \Phi)|^2 [f_n^c(k_{||}, \Phi) - f_m^v(k_{||}, \Phi)] L(\omega, k_{||}, \Phi) = \frac{1}{2} \int_0^{2\pi} d\Phi \int_0^{2\pi} d\Phi \int_0^{2\pi} dk_{||} \frac{2k_{||}}{(2\pi)^2 L_w} |M_{nm}(k_{||}, \Phi)|^2 [f_n^c(k_{||}, \Phi) - f_m^v(k_{||}, \Phi)] L(\omega, k_{||}, \Phi) = \frac{1}{2} \int_0^{2\pi} d\Phi \int_0^{2\pi} d\Phi \int_0^{2\pi} dk_{||} \frac{2k_{||}}{(2\pi)^2 L_w} |M_{nm}(k_{||}, \Phi)|^2 [f_n^c(k_{||}, \Phi) - f_m^v(k_{||}, \Phi)] L(\omega, k_{||}, \Phi)$$

$$L(\omega, k_{||}, \Phi) = \frac{(1 - \operatorname{Re}Q(k_{||}, \hbar\omega))\operatorname{Re}\Xi(E_{lm}(k_{||}, \hbar\omega)) - \operatorname{Im}Q(k_{||}, \hbar\omega)\operatorname{Im}\Xi(E_{lm}(k_{||}, \hbar\omega))}{(1 - \operatorname{Re}Q(k_{||}, \hbar\omega))^2 + (\operatorname{Im}Q(k_{||}, \hbar\omega))^2}$$

$$\begin{aligned} \operatorname{Re}\Xi(E_{lm}(k_{||},\hbar\omega)) &= \sqrt{\frac{\pi\tau_{co}}{2\hbar\Gamma_{cv}(k_{||},\hbar\omega)}} \\ &\times \exp\left(-\frac{\tau_{co}}{2\hbar\Gamma_{cv}(k_{||},\hbar\omega)}E_{lm}^2(k_{||},\hbar\omega)\right) \end{aligned}$$

$$\begin{split} \mathrm{Im}\Xi(E_{lm}(k_{||},\hbar\omega)) &= \frac{\tau_{co}}{\hbar} \int_{0}^{\infty} \exp\left(-\frac{\Gamma_{cv}(k_{||},\hbar\omega)\tau_{co}}{2\hbar}t^{2}\right) \\ &\times \sin\left(\frac{\tau_{co}E_{lm}(k_{||},\hbar\omega)}{\hbar}t\right) dt. \end{split}$$

|M|<sup>2</sup>: Momentum matrix element

 $\boldsymbol{f}_n$  and  $\boldsymbol{f}_m$  : Fermi functions

L : Gaussian line shape function renormalized with many-body effects

Q: CE many-body effect

 $\Phi$  : Angle between kx' and ky' wavevectors



### InGaN/GaN QW Structure





### InGaNAs/GaAs QW structure





### InGaAsP/InP QW structure



IEEE JQE 35, 771(1999)



### Lattice constant for MgZnO



 <sup>3</sup>T. Makino, Y. Segawa, M. Kawasaki, A. Ohtomo, R. Shiroki, K. Tamura, T. Yasuda, and H. Koinuma, Appl. Phys. Lett. **78**, 1237 (2001).
 <sup>4</sup>A. Ohtomo and A. Tsukazaki, Semicond. Sci. Technol. **20**, S1 (2005).

 $P_{SP}^{\,\rm ZnO}\!=-\,0.05 C/m^2$ 

Spontaneous polarization constant for MgO : a fitting parameter



# Energy shift and exciton energy for ZnO/MgZnO QW structures





# Difference between total PZ and SP polarizations and potential energy





# Optical gain spectra and optical matrix elements





# Polarization and Internal field of ZnO/MgZnO QW structures





### Potential profile : ZnO/MgZnO





### Potential profile : InGaN/GaN





### Valence band structure: InGaN/GaN





### Valence band structure: ZnO/MgZnO





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### Average hole effective mass





### Quasi-Fermi level separation





### Optical matrix elements





### Optical gain



➢In the case of small crystal angle, ZnO system has much larger optical gain than the GaN systems. This is because WZ ZnO/MgZnO QW structure has the larger matrix element and smaller effective mass than GaN-based QW structures near (0001) crystal orientation.

➢On the other hand, in the case of the (1010) crystal orientation, InGaN/GaN QW structures show larger optical gain than ZnO/MgZnO QW structures due to the larger matrix element and the smaller effective mass.

GaN/AlGaN : Crystal orientation effect is relatively small.



### Summary

- The negligible internal field observed in the ZnO/MgZnO QW structure with small Mg compositions and thin well widths can be explained by the cancellation of total PZ and SP polarizations in the well and that in the barrier.
- In the case of (0001) crystal orientation, ZnO/MgZnO QW structure has much larger optical gain than the GaN-based QW structure.
- On the other hand, in the case of the (1010) crystal orientation, InGaN/GaN QW structure has larger optical gain than ZnO/MgZnO QW structures.
- GaN/AIGaN : Crystal orientation effect is relatively small.



### Appendix



### Line shape function





### Constant energy contour



InGaN/GaN QW structure

#### ZnO/MgZnO QW structure

